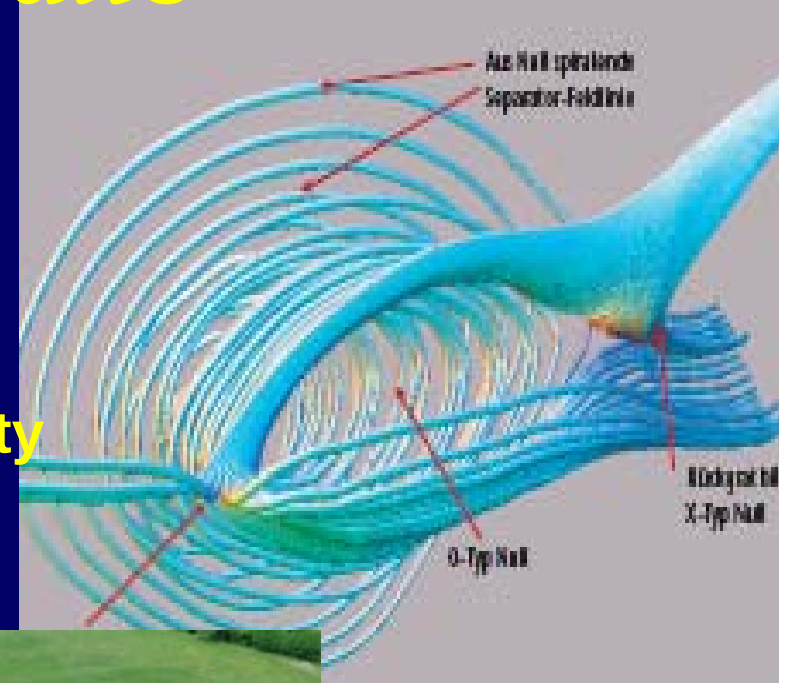


What is driving the Solar atmosphere?

Magnetic reconnection, RHMD and kinetic simulation results



Jörg Büchner
MPI for Solar
System Research
Katlenburg-Lindau
and Göttingen University
in Germany



The Sun as a star

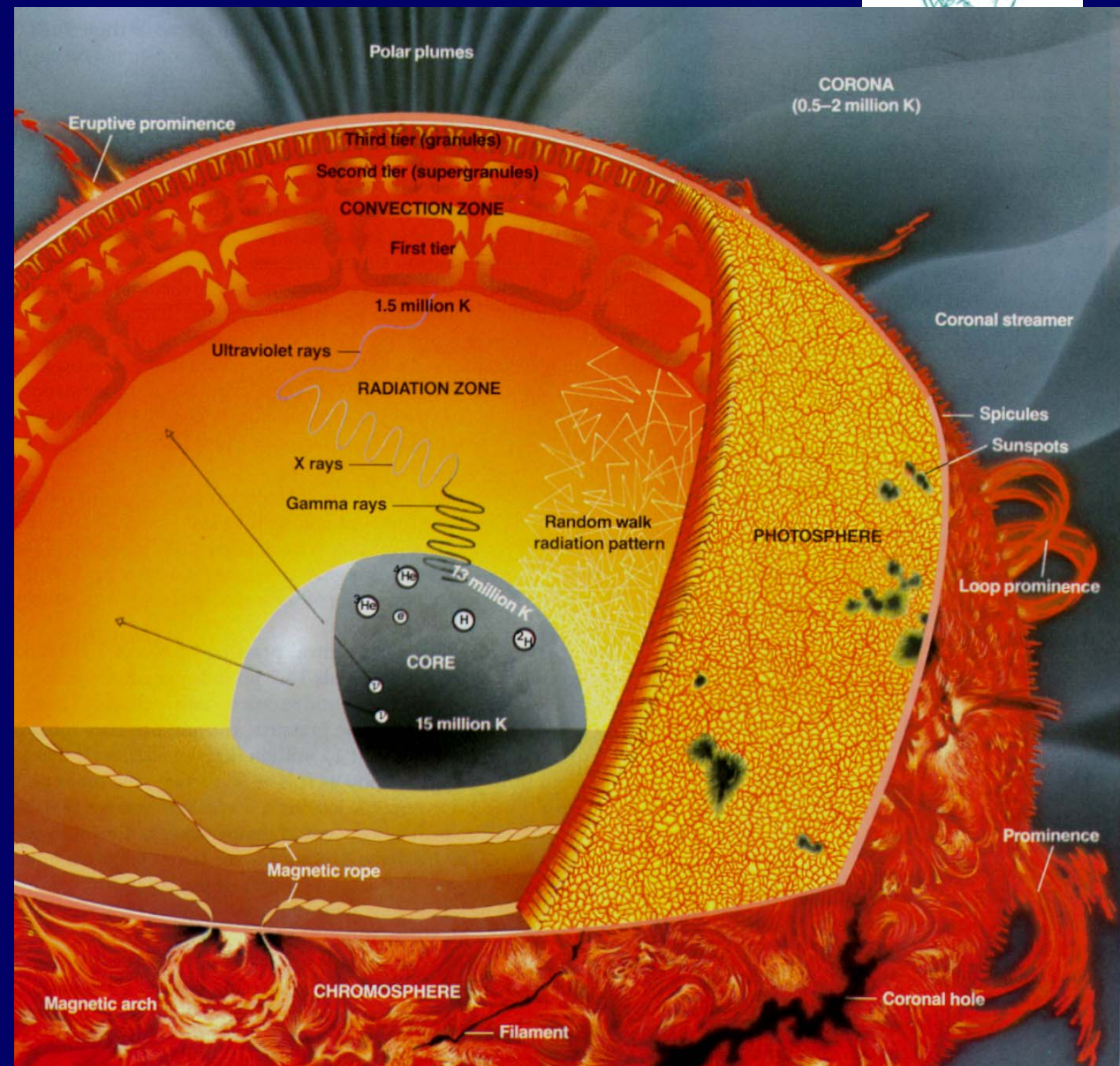


Parent galaxy	Milky Way
Type	fixed star
Spectral class	G2
Magnitude	+ 4.8
Distance to Earth	149 598 000 km i.e., 1 AU
Radius R_s	696,000 km i.e., 109 R_E
Total mass M_s	1.989×10^{30} kg i.e., 333 000 M_E
Density (average)	1.409 g cm^{-3}
Surface temperature	5800 K
Rotation duration	25 days at equator. 35 days at poles.
Age	4.60 billion years
Number of planets	8 by definition
Next neighbor star	Alpha-Centauri, at 4.37 light years
Next neighbor galaxy	Magellan's Clouds, at 165.000 light years
Earth's distance variation	+/- 1.69 % (+ in July, - in January)
Apparent diameter	31' 59.3" = 1913.3 " i.e. 0.5 degree
Apparent radius	959.65" i.e. 1000 arcsec
1 arcsec on sun, from Earth	725 km
Energy output	3.82×10^{33} Watt
Energy input into Earth	1.370 Watt/m^2
total	173 Mio Gigawatt

The Sun as an onion



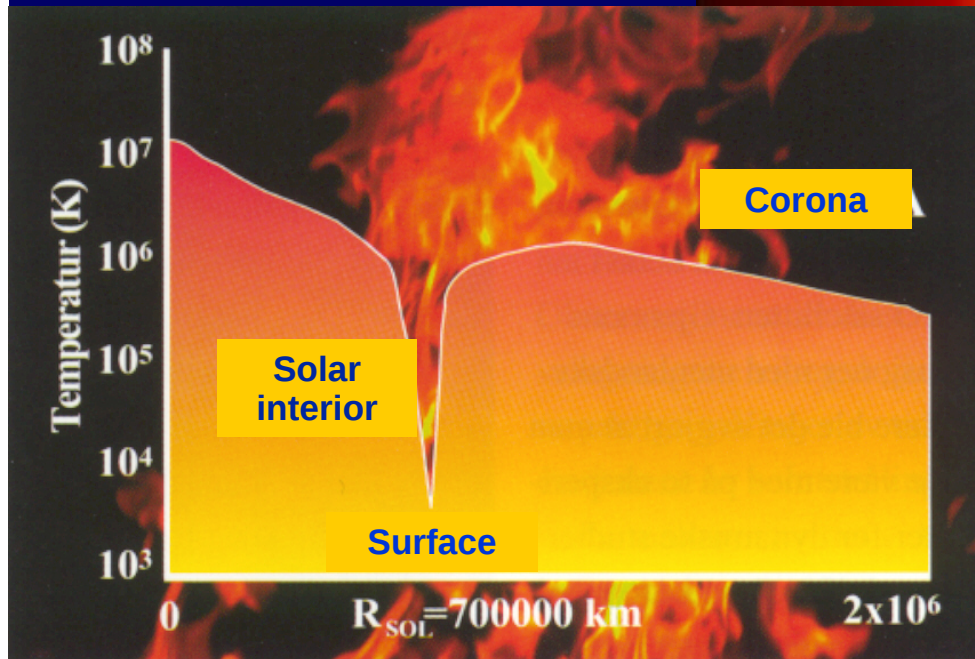
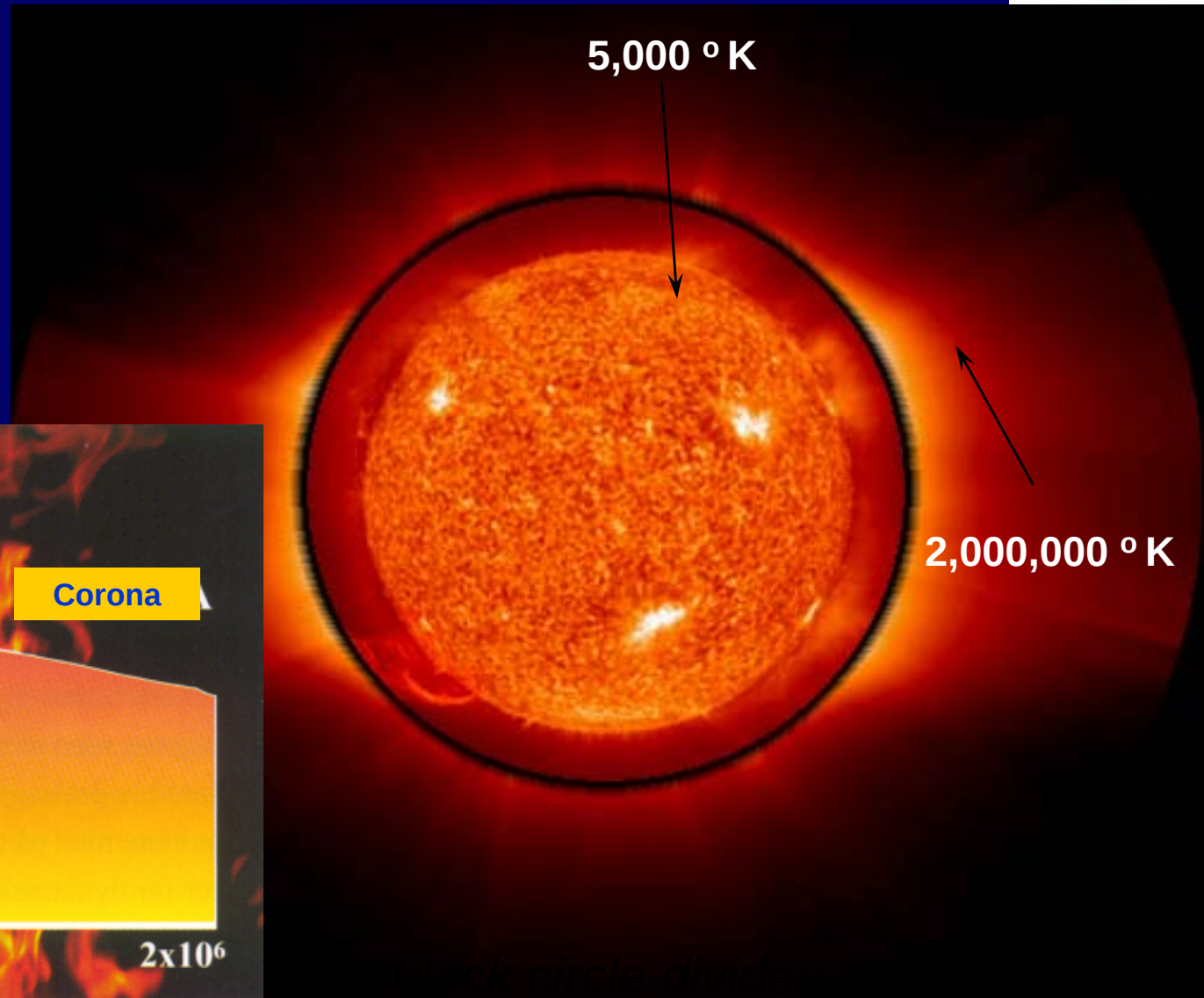
- Source in the **Core**: 700 Million tons of He burned per second
- **Radiation zone**: Photon transport over Millions of years
- **Convection zone**: Cellular convective transport
- **Photosphere**: Plasma cools down to 6000 K
- **Atmosphere: Chromosphere & Corona**:
 - Temperature rises to millions of K
 - Solar Wind ejected



The atmosphere of the Sun



Above the photosphere the temperature rises from 6000 K to millions of K and the hot plasma is accelerated into the solar wind



Dynamics: 150 years of flare observations

MONTHLY NOTICES

OF THE

ROYAL ASTRONOMICAL SOCIETY.

VOL. XX.

Nov. 11, 1859.

No. I.

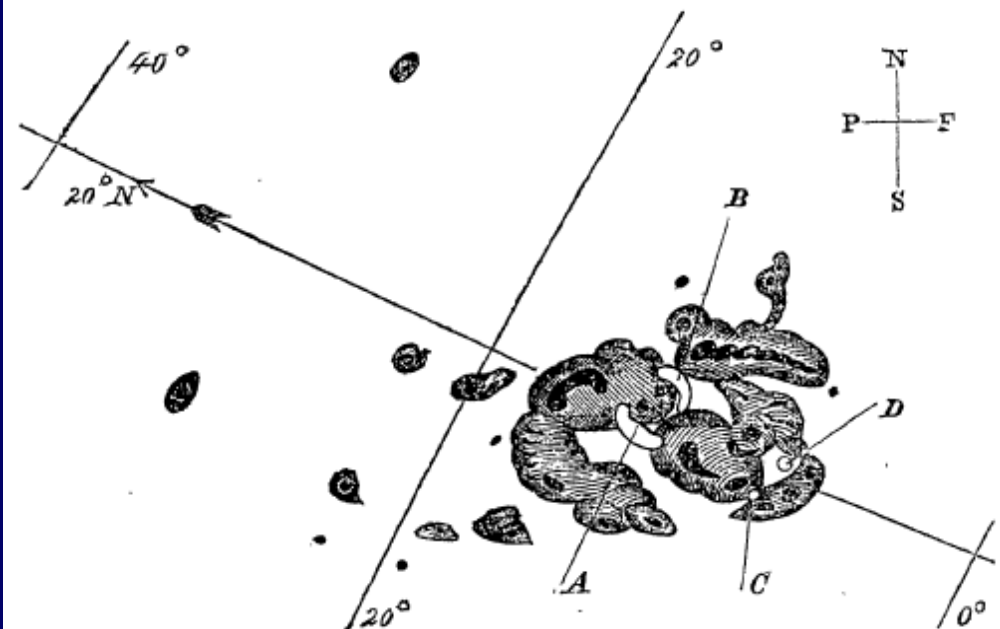
It has been very gratifying to me to learn that our friend Mr. Hodgson chanced to be observing the sun at his house at Highgate on the same day, and to hear that he was a witness of what he also considered a very remarkable phenomenon. I have carefully avoided exchanging any information with that gentleman, that any value which the accounts may possess may be increased by their entire independence.

(Mr. Carrington exhibited at the November Meeting of the Society a complete diagram of the disk of the sun at the time, and copies of the photographic records of the variations of the three magnetic elements, as obtained at Kew, and pointed out that a moderate but very marked disturbance took place at about 11^h 20^m A.M., Sept. 1st, of short duration; and that towards four hours after midnight there commenced a great magnetic storm, which subsequent accounts established to have been as considerable in the southern as in the northern hemisphere. While the contemporary occurrence may deserve noting, he would not have it supposed that he even leans towards hastily connecting them. "One swallow does not make a summer.")

Description of a Singular Appearance seen in the Sun on September 1, 1859. By R. C. Carrington, Esq.

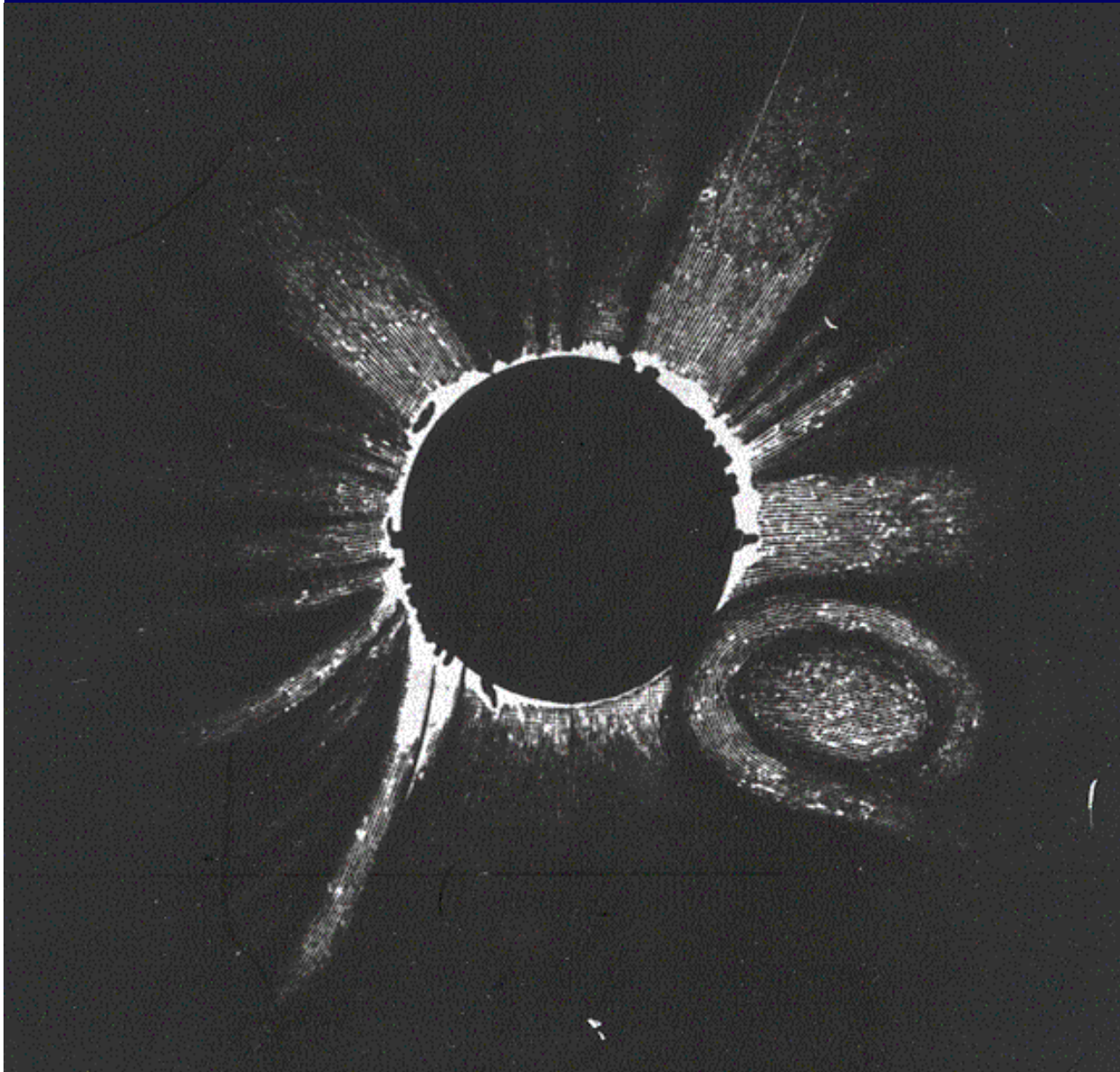
While engaged in the forenoon of Thursday, Sept. 1, in taking my customary observation of the forms and positions of the solar spots, an appearance was witnessed which I believe to be exceedingly rare. The image of the sun's disk was, as usual with me, projected on to a plate of glass coated with distemper of a pale straw colour, and at a distance and under a power which presented a picture of about 11 inches diameter.

I had secured diagrams of all the groups and detached spots, and was engaged at the time in counting from a chronometer and recording the contacts of the spots with the cross-wires used in the observation, when within the area of the great north group (the size of which had previously excited general remark), two patches of intensely bright and white light broke out, in the positions indicated in the appended diagram by the letters A and B, and of the forms of the spaces left white. My



first impression was that by some chance a ray of light had penetrated a hole in the screen attached to the object-glass, by

... and plasma detachments



Maybe, the oldest drawing of a plasma detachment from the Sun (or what is now called „CME“, a Coronal Mass Ejection) seen 1860 in Tempel, Spain [Ranyard, 1879]

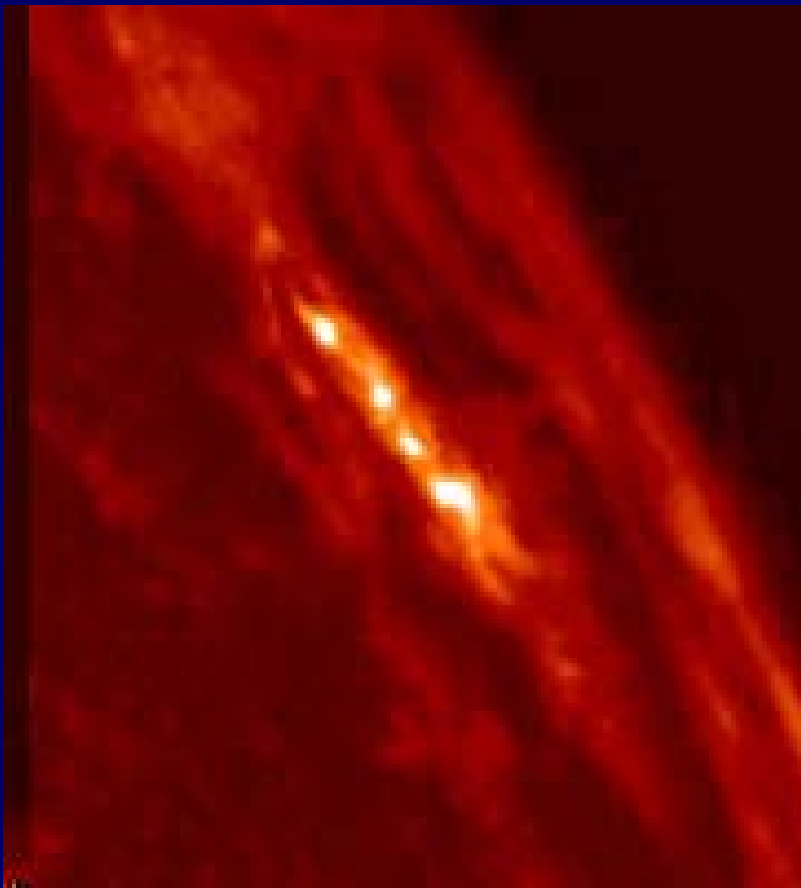
Flares -> Quakes



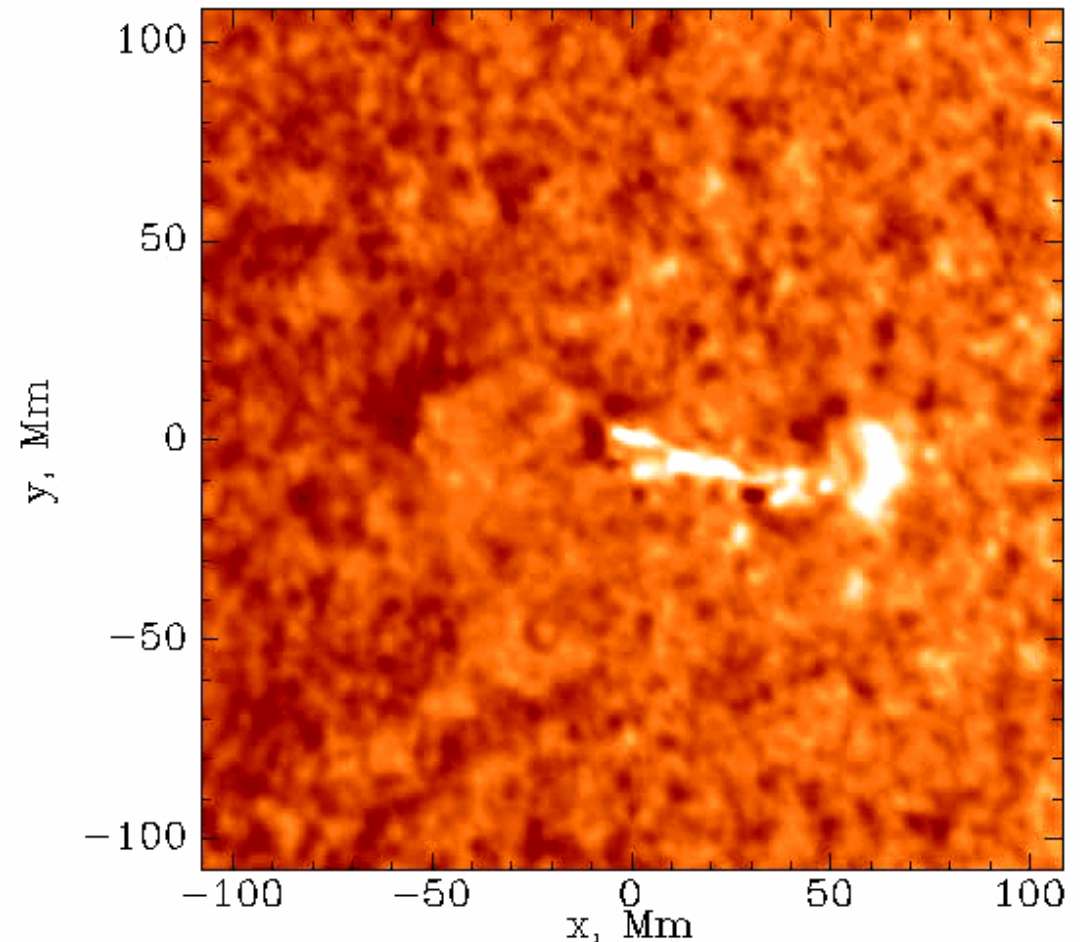
Nowadays: direct observations Flares as sudden eruptions of radiation, especially from so called „active regions“ (sunspots)

Flares -> Quakes

TRACE s/c



SOHO-MDI



Helioseismology of the interior

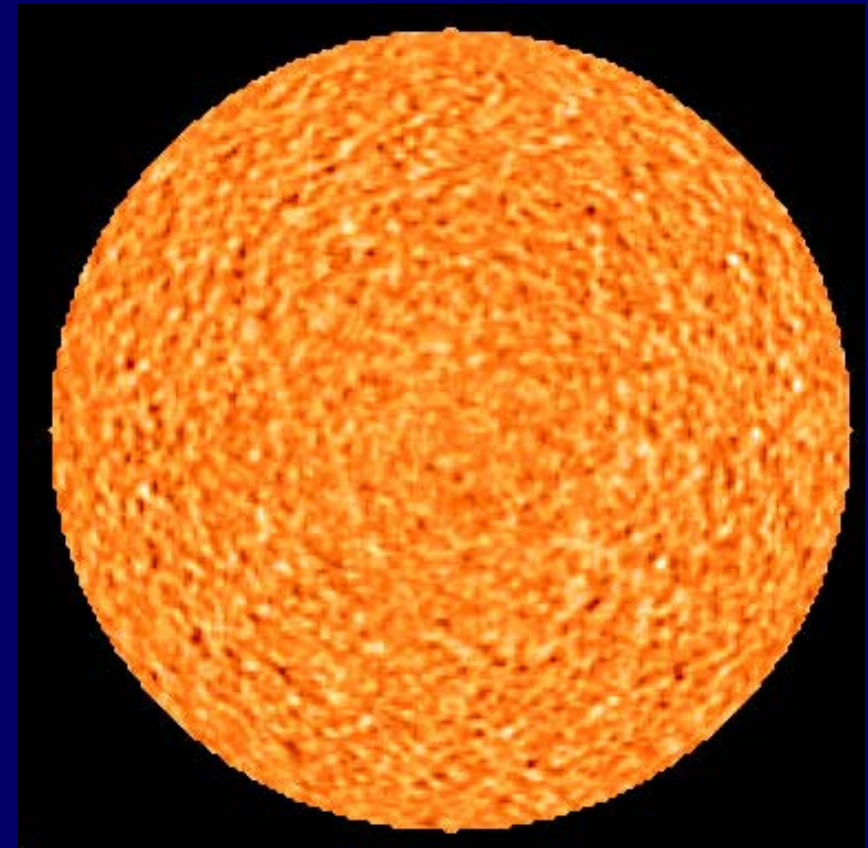
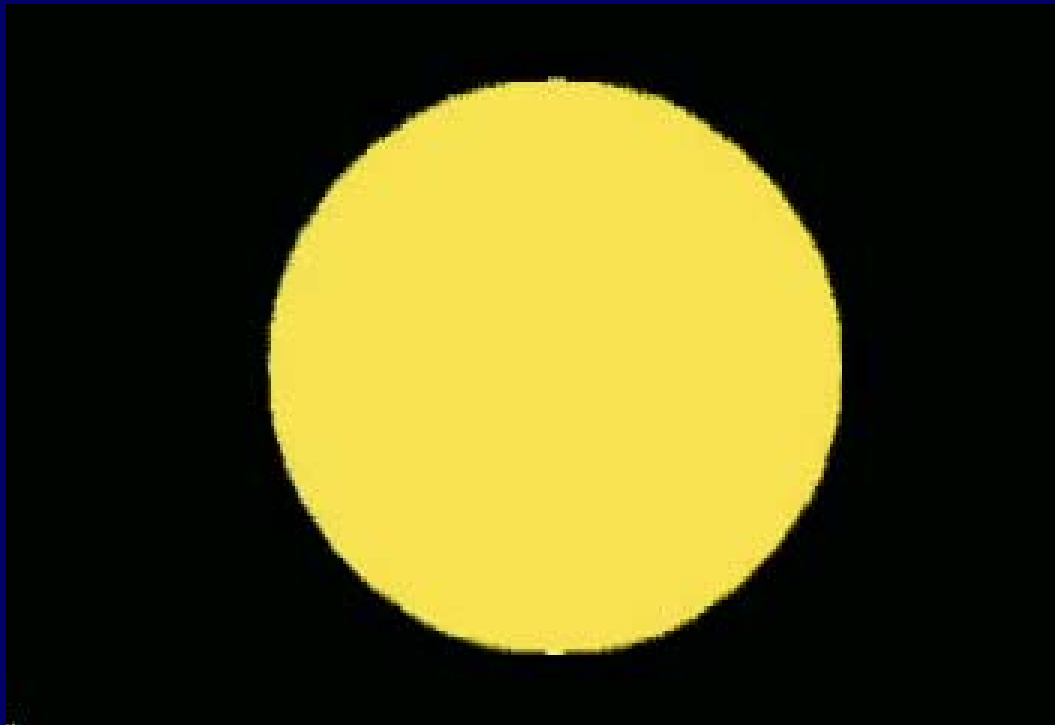


- Quakes and vibrations of the sun are used for diagnostics:
- The internal wave reflections tell us more about the interior of the sun
-> helioseismology:

Acoustic signal:



(42 000 times accelerated)

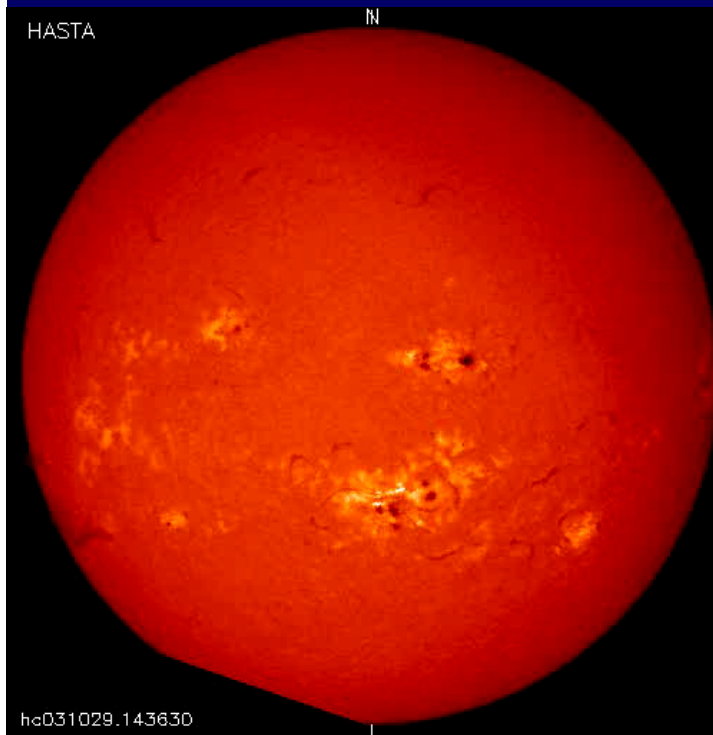


Observations of the atmosphere



Observation of two strong Flares:
28.10.2003 (X17.2) and
29.10.2003 (X10.0)

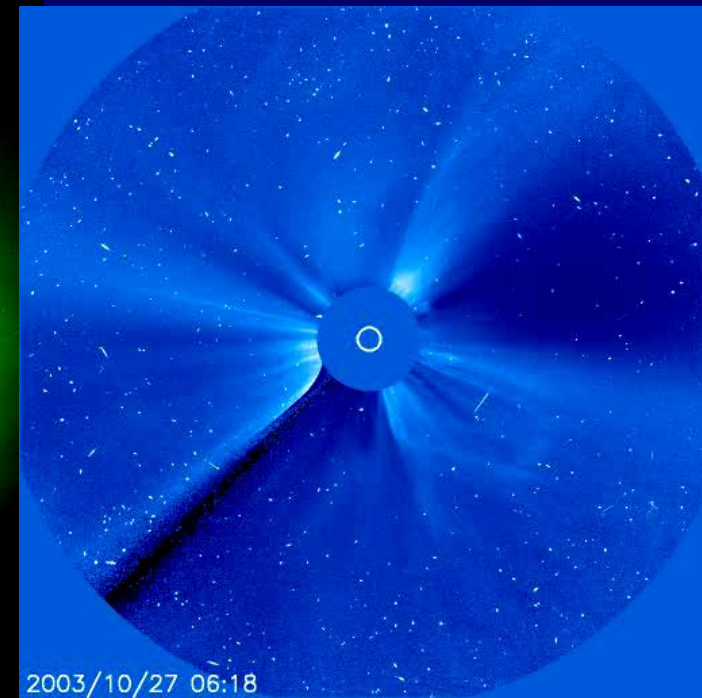
(„snow storm“: due to
energetic particles that
disturb the camera
CCD)



**B-field (HASTA,
El Leoncito)**

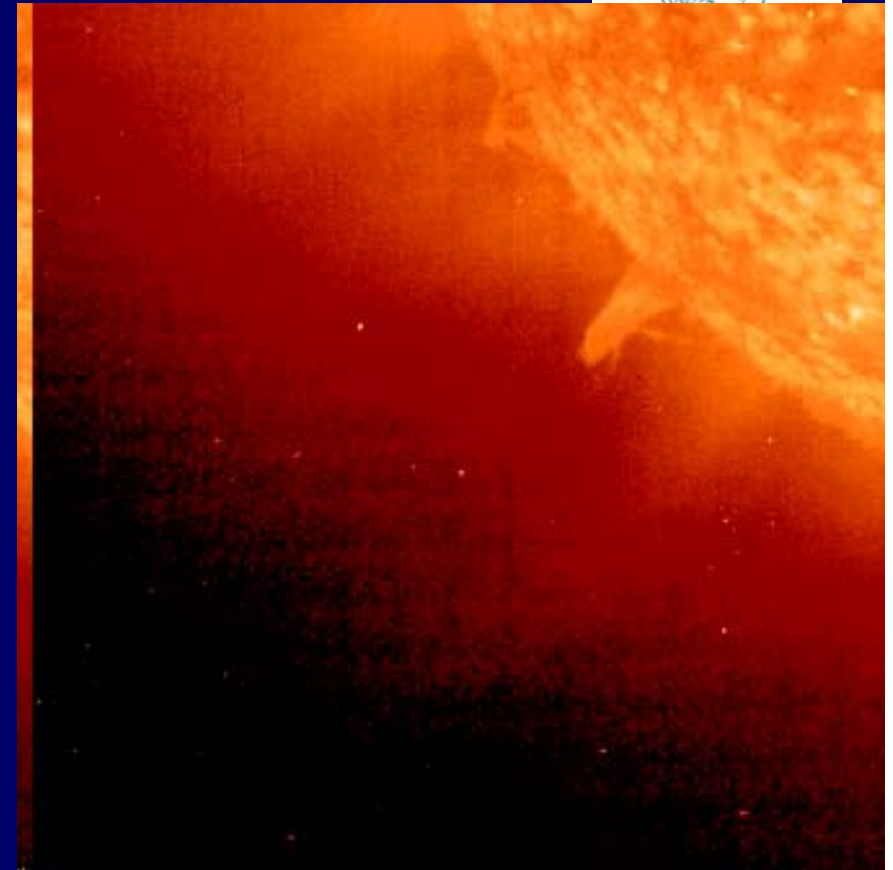
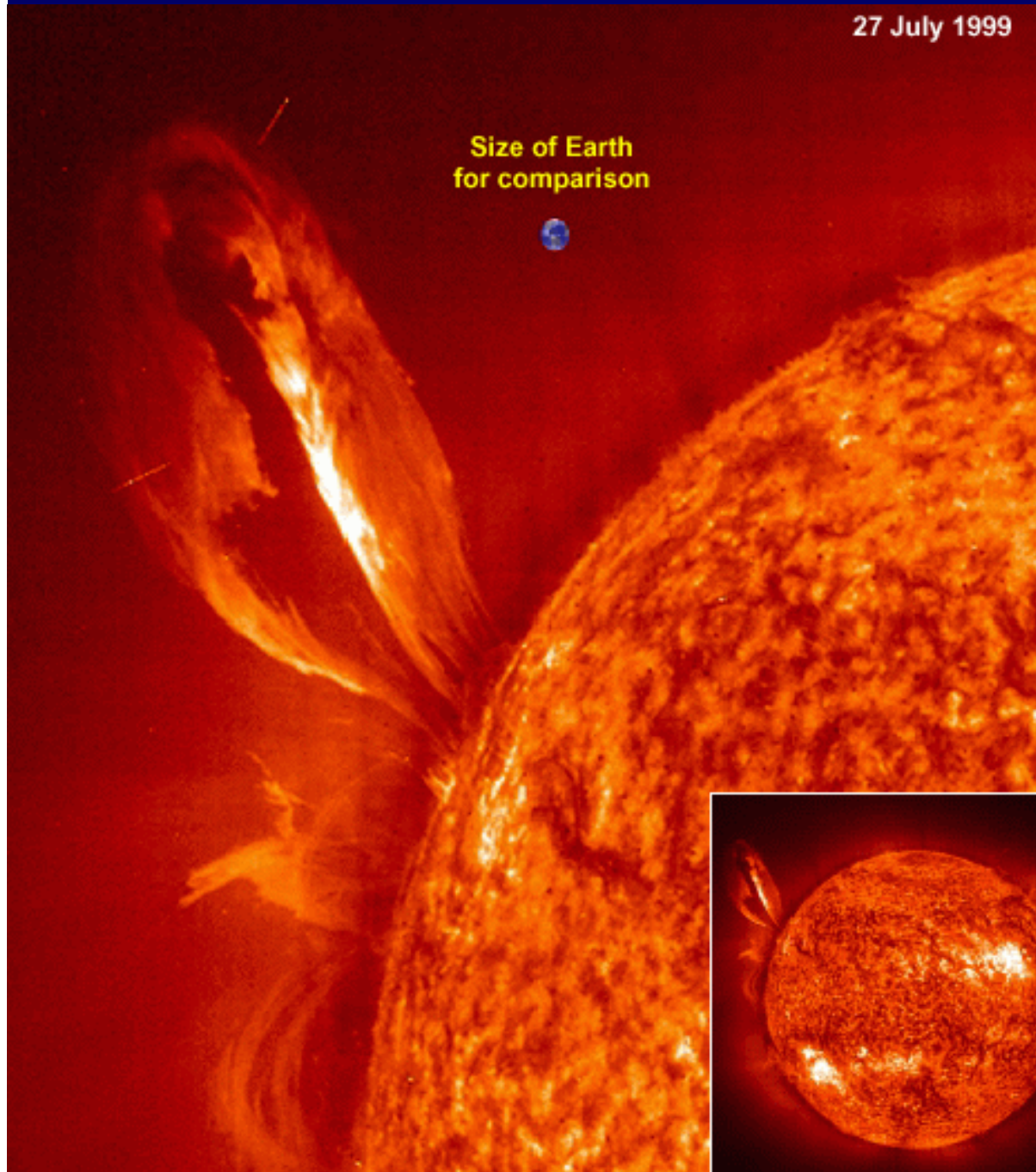


**EUV Image
(EIT on SOHO)**



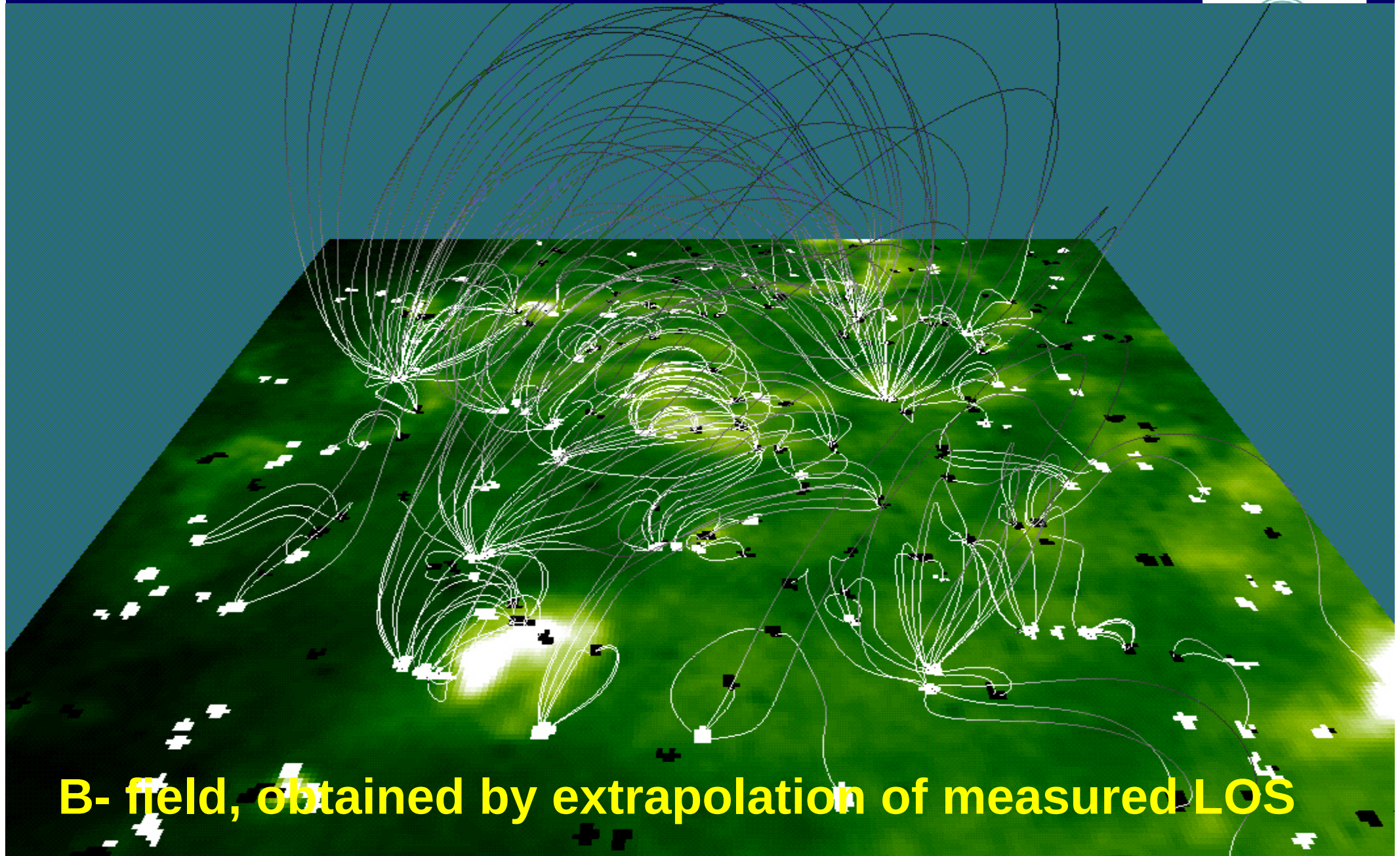
**Coronagraphy
(LASCO-C3, SOHO)**

Prominence eruptions



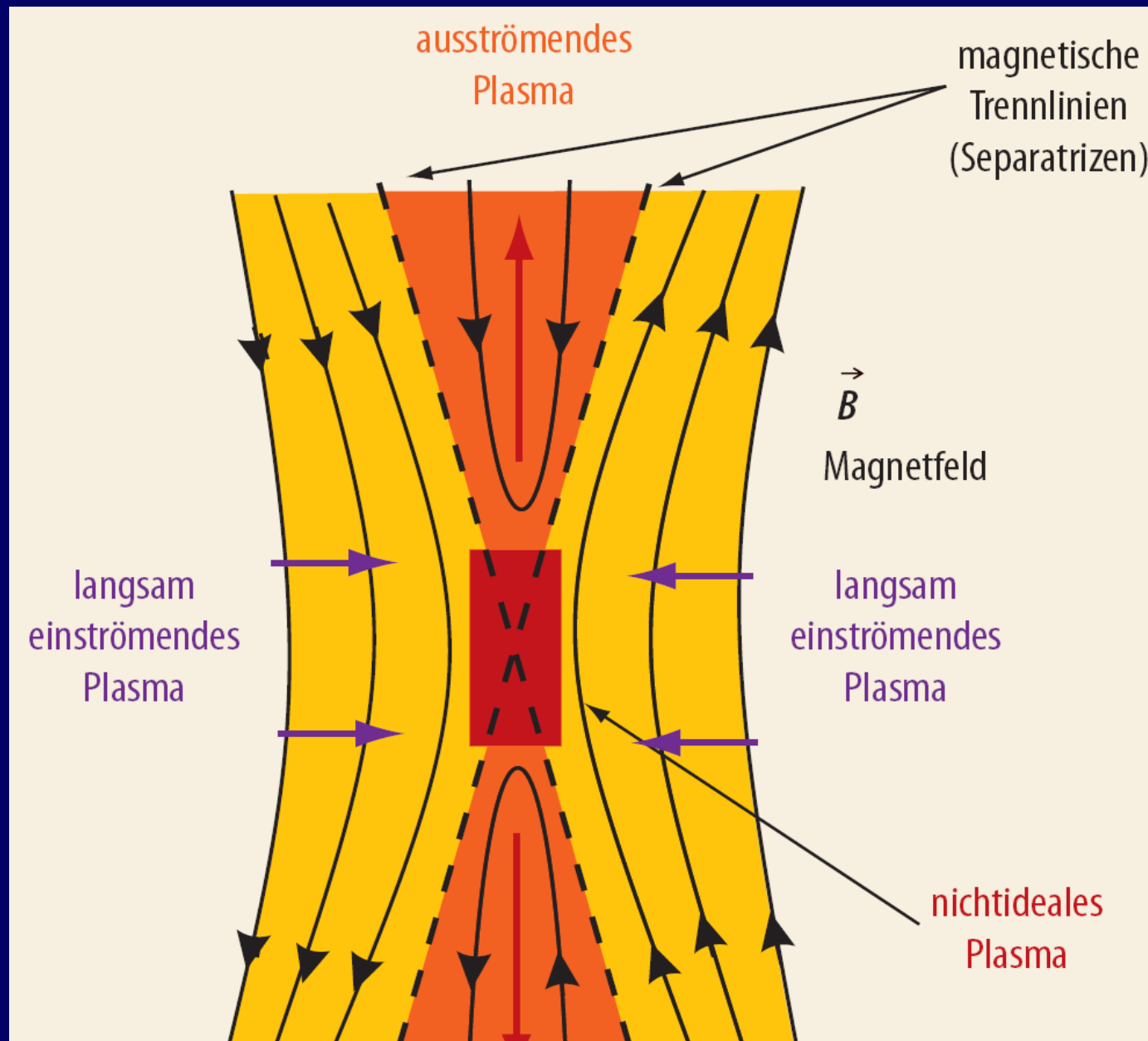
**Prominence
eruptions
-> CMEs**

Extrapolation: magnetic carpet



B- field, obtained by extrapolation of measured LOS

Energy release by reconnection

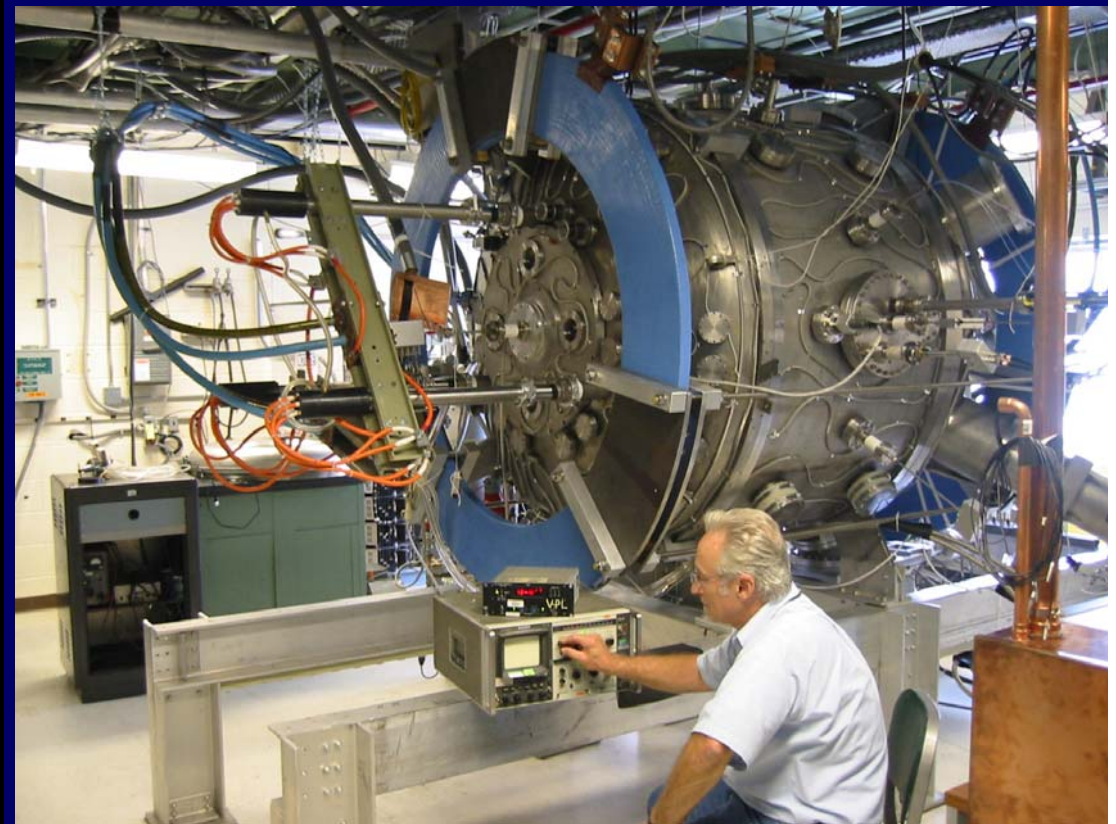
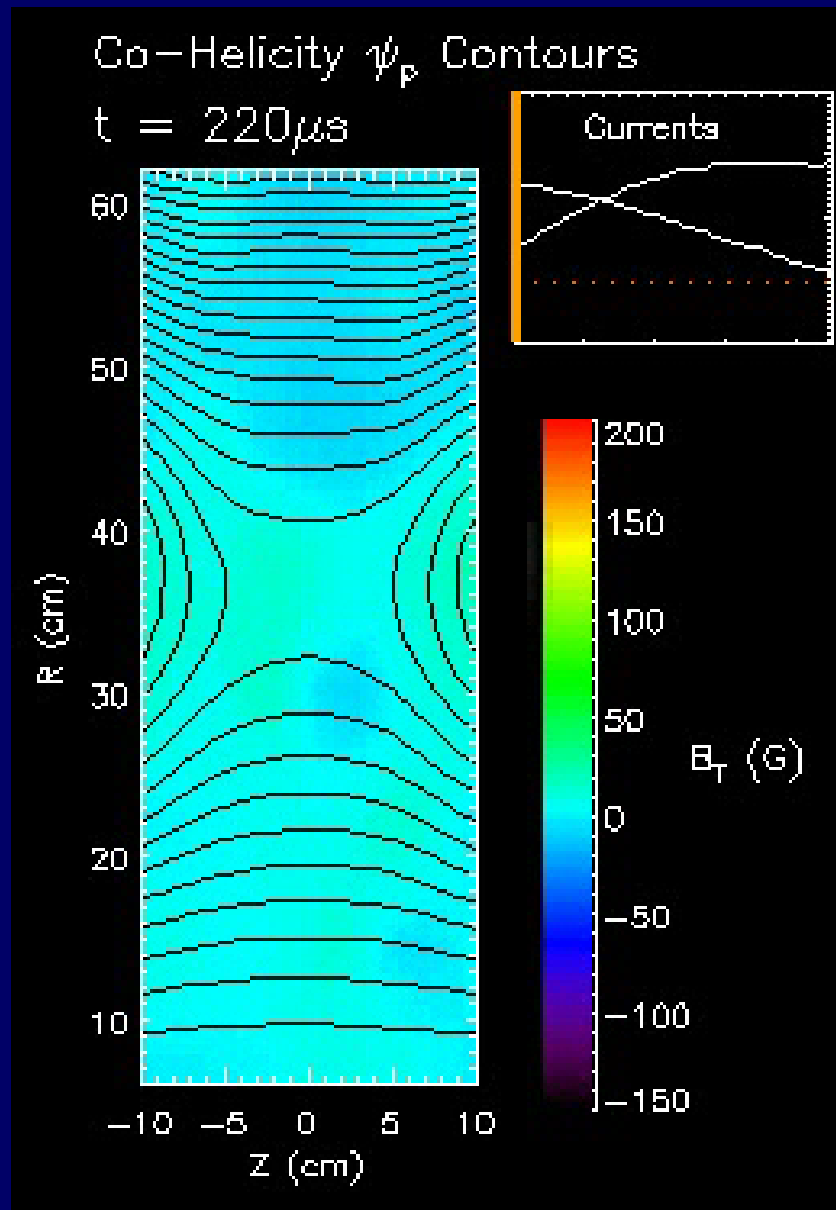


**Demonstration
of the principle
action of
reconnection in
two dimensions**

Reconnection in the laboratory



Laboratory-experiment by
Yamada et al.: MRX
in the PPP Lab Princeton



Elements to be considered in modelling the solar atmosphere:



- Energy input and B-fields from the photosphere
 - Magnetic energy release in the solar atmosphere
 - Plasma and neutral gas stratification in the solar atmosphere by gravity
 - In the chromosphere: coupling to neutral gas
 - What is causing energy dissipation in the atmosphere?
 - Important: heat conduction, radiative losses
- > Appropriate model - description and applications
- Solar wind acceleration
 - Heating of Bright Points

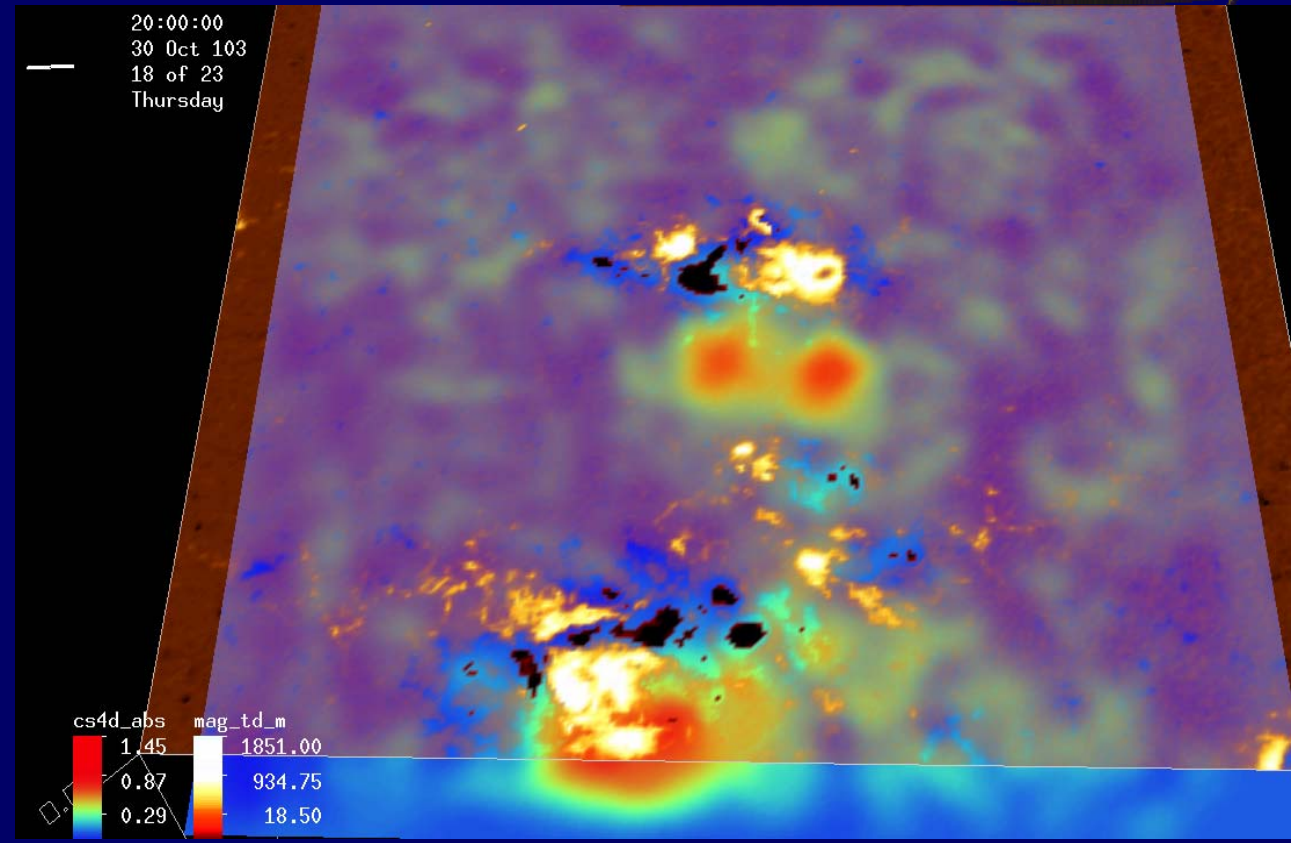
Energy source for the corona:

MPS

Plasma convection
below the photosphere
(to the right: helioseis-
mology of AR10488,
30.10.03, lower panel:
16 Mm deep) [Gizon &
Kosovichev]

-> Dynamo -> B fields

-> upward Poynting flux
estimated for:



$$E = -v \times B; v = 0.1 \text{ km} \cdot \text{s}^{-1}; B = 100 \text{ G}$$

Vs. observed fluxes:

Quiet regions 300 W m^{-2}

Active regions $(0.5 - 1) 10^4 \text{ W m}^{-2}$

Coronal holes 800 W m^{-2}

$$\frac{\mathbf{E} \times \mathbf{B}}{\mu} \approx \frac{v B^2}{\mu} \approx 10^4 \text{ W} \cdot \text{m}^2$$

Large scale modelling: RHMD

MPS



$$\begin{aligned}\frac{\partial \rho}{\partial t} &= -\vec{\nabla} \cdot \rho \vec{u} - \rho (\vec{u} \cdot \vec{\nabla}) \vec{u} - \rho_0 \\ \frac{\partial \rho \vec{u}}{\partial t} &= -\vec{\nabla} \cdot \rho \vec{u} \vec{u} - \vec{\nabla} p + \vec{j} \times \vec{B} - \nu \rho (\vec{u} - \vec{u}_0) \\ &= -\vec{\nabla} \cdot \left[\rho \vec{u} \vec{u} + \left(p + \frac{B^2}{2\mu_0} \right) \mathbf{1} - \frac{\vec{B} \vec{B}}{\mu_0} \right] - \nu \rho (\vec{u} - \vec{u}_0) \\ \frac{\partial \vec{B}}{\partial t} &= \vec{\nabla} \times (\vec{u} \times \vec{B} - \eta \vec{j}) \\ \frac{\partial p}{\partial t} &= \end{aligned}$$

Closing energy equation

The index ₀ indicates density, temperature and velocity of the neutral gas component, which in chromosphere and photosphere is strongly coupled to the plasma!

$$p = 2n\kappa_B T$$

$$\begin{aligned}\vec{E} &= -\vec{u} \times \vec{B} + \eta \vec{j} \\ \vec{\nabla} \times \vec{B} &= \mu_0 \vec{j}\end{aligned}$$

$$\hat{\eta}^* = \begin{cases} \eta_a^* \left(\frac{|v_{dr}|}{v_{thr}} - 1 \right); & |v_{dr}| \geq v_{thr}, \\ 0; & |v_{dr}| < v_{thr}. \end{cases}$$

Energy equation

$$\frac{\partial p}{\partial t} = -\nabla \cdot pu - (\gamma - 1)p\nabla \cdot u + (\gamma - 1)S$$

where

heat
conduction

along
magnetic
field

$$\nabla \cdot q = \nabla_{\parallel} \cdot (\kappa_{\parallel} \nabla_{\parallel} T)$$

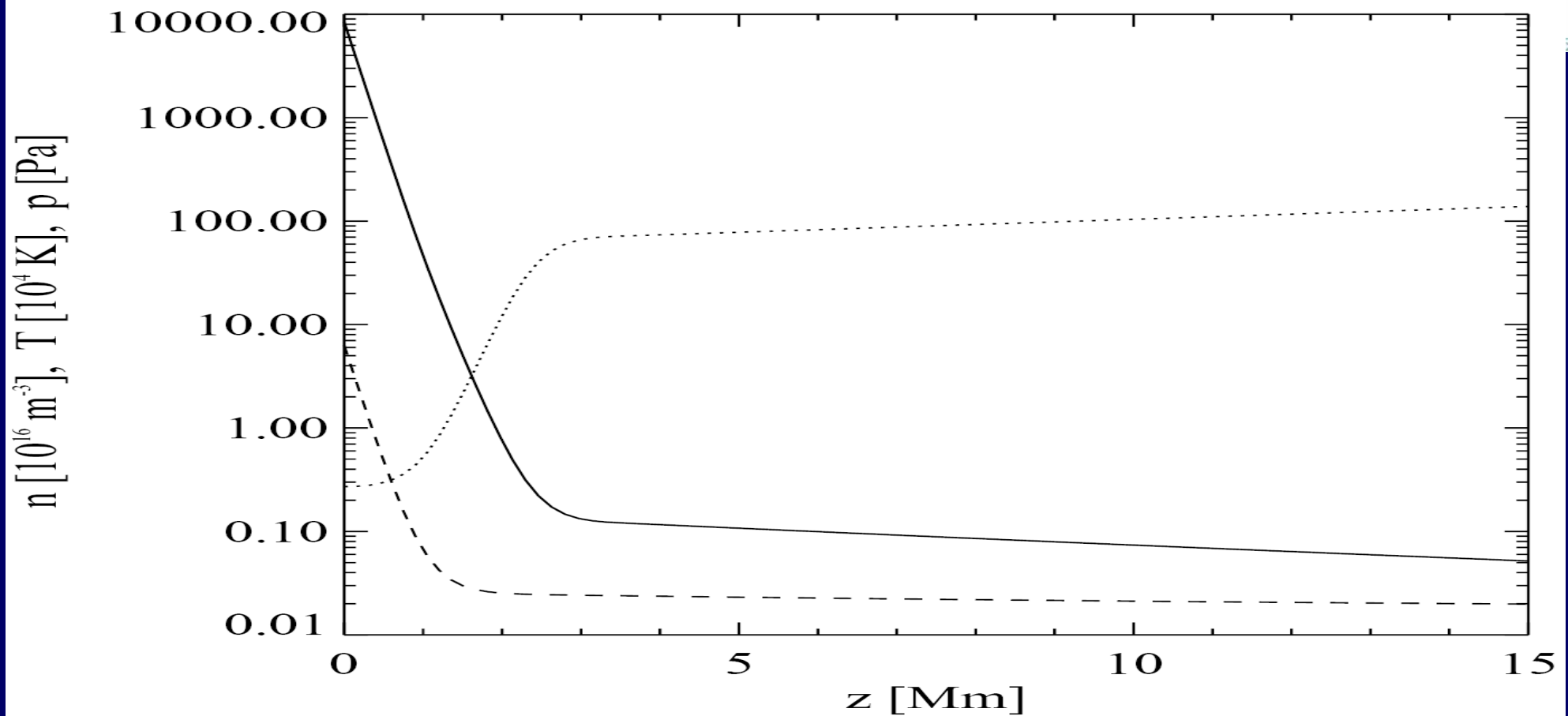
$$S = \eta j^2 - \nabla \cdot q - L_r$$

$$L_r = n n_n \chi T^{\alpha} \text{ W/m}^3$$

ohmic dissipation

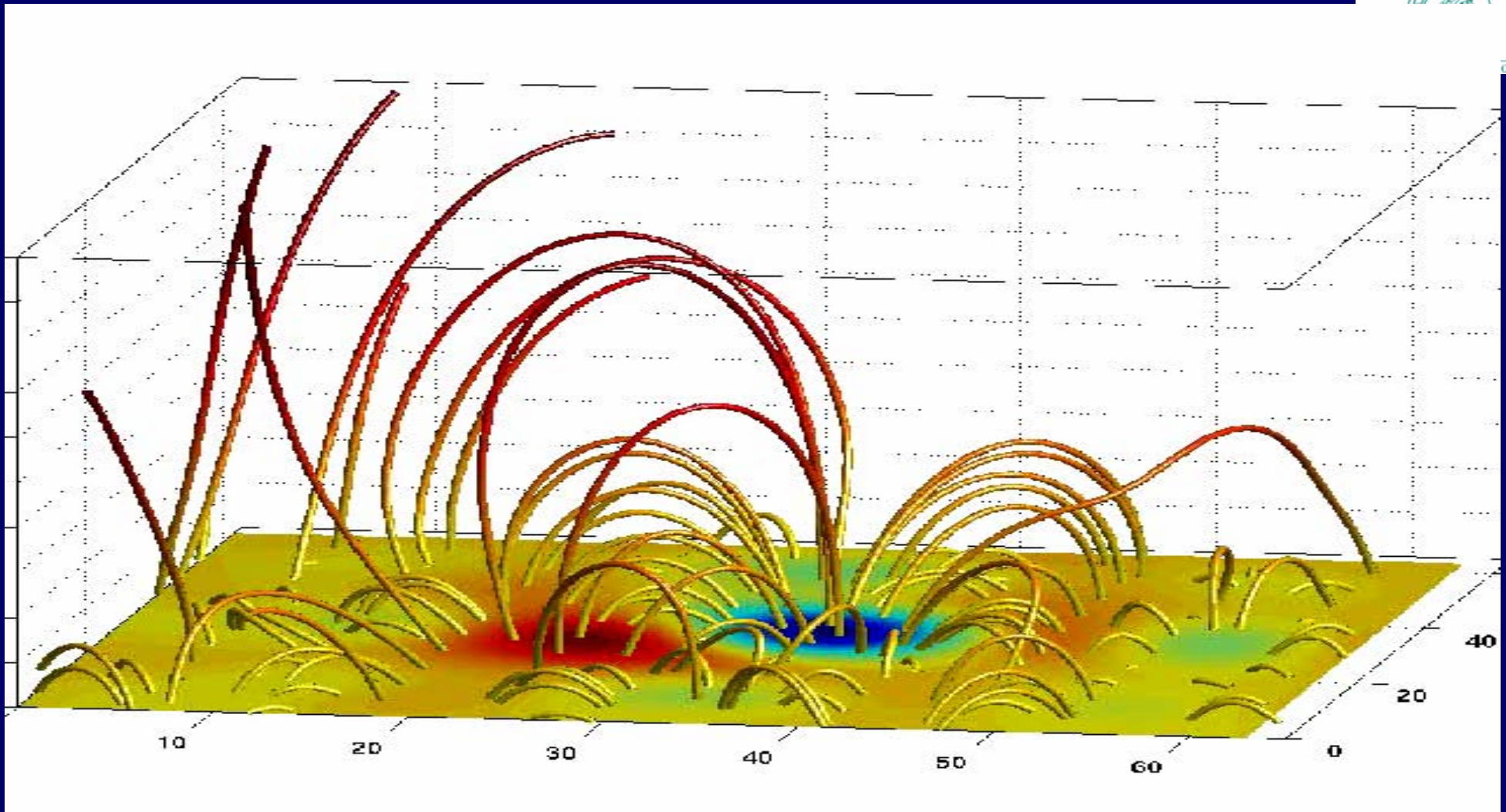
radiative cooling

Initial inhomogeneity



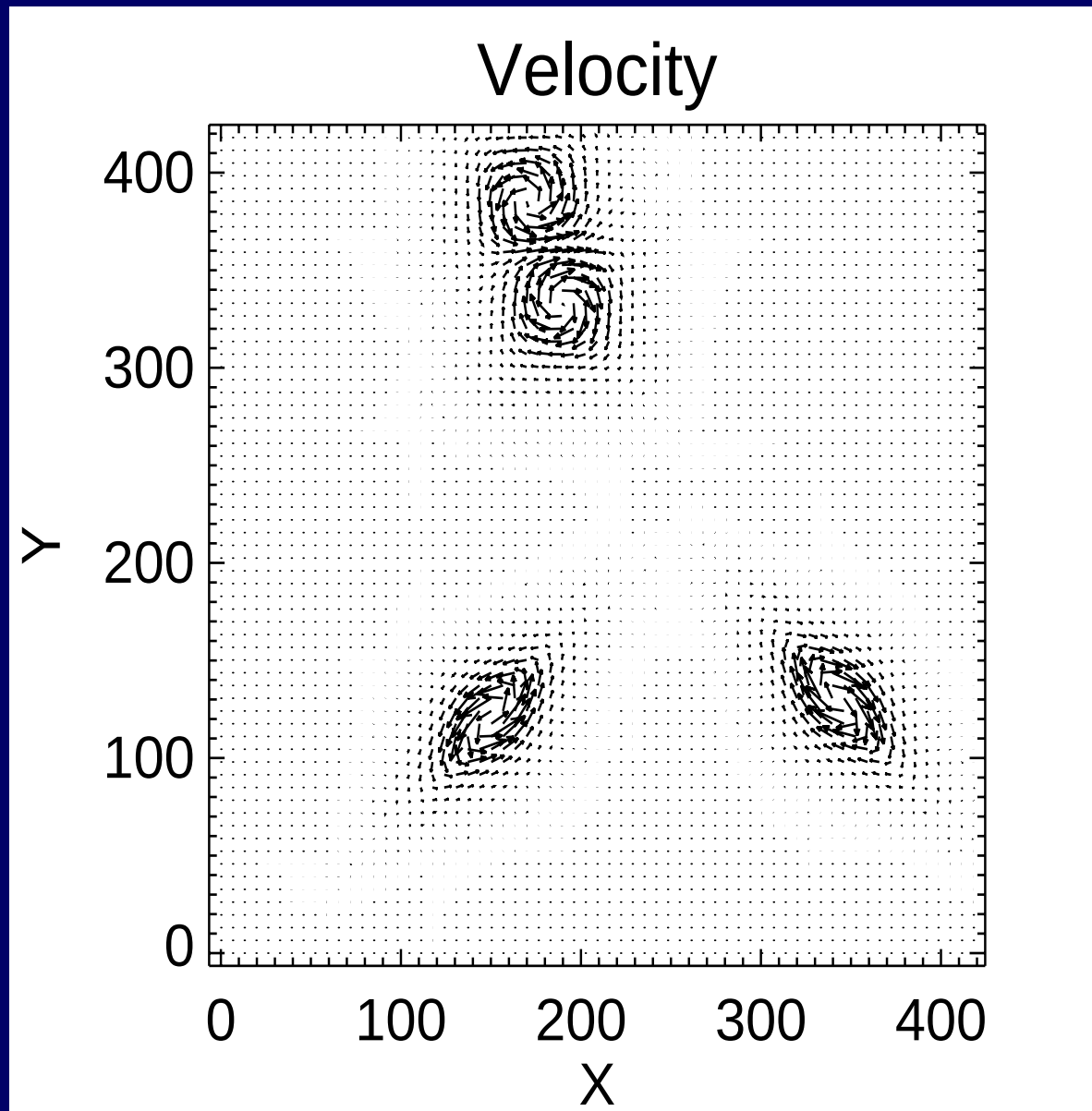
Plasma density, temperature and pressure are radially stratified according to the solar gravitation

Extrapolated initial B-field



The initial magnetic field is extrapolated from the observed photospheric magnetic field

Observed plasma motion as a boundary condition



By correlating the line-of sight B-field patterns one can derive the approximate motion picture as shown and use it as a boundary condition for the simulation. The motion creates currents in the corona

-> current dissipation?

-> magnetic reconnection?

Dissipation in resistive MHD



Induction equation

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}$$

... reveals scales -> magnetic Reynolds number

$$R_m = \frac{\mu_0 l v}{\eta}$$

**Rm must be small, hence,
in the solar atmosphere**

1.) $l \sim 1 \text{ km}$

2.) $\eta \sim 10^8$ times that of collisions

How can these scales / dissipation be reached?

Dissipation in multi-fluids



A two-fluid-description reveals the „generalized Ohm's law“ of magnetohydrodynamics (MHD):

$$\frac{4\pi}{\omega_{pe}^2} \frac{d\vec{J}}{dt} = \boxed{\vec{E} + \vec{v}_i \times \vec{B}} - \frac{1}{ne} \vec{J} \times \vec{B} + \frac{1}{ne} \nabla p_e - \eta \vec{J}$$

c/ω_{pe}

<- spatial ->
<- scales ->

ρ_i

c/ω_{pi}

$\wedge$

electron
inertia

elektrons
- ions
decoupled,
„Hall“ term“

off-diagonal
elements
of the
pressure
tensor

dissipation
due to
turbulence

Microscale (kinetic) dissipation



Ensemble averaging:

$$\langle \delta f_j \rangle = \langle \delta \vec{E} \rangle = \langle \delta \vec{B} \rangle = 0.$$

$$f_j = f_{0j} + \delta f_j \quad E_{\parallel} = \langle E_{\parallel} \rangle + \delta E_{\parallel}$$

-> Modified Vlasov equation, after velocity averaging -> momentum exchange rate

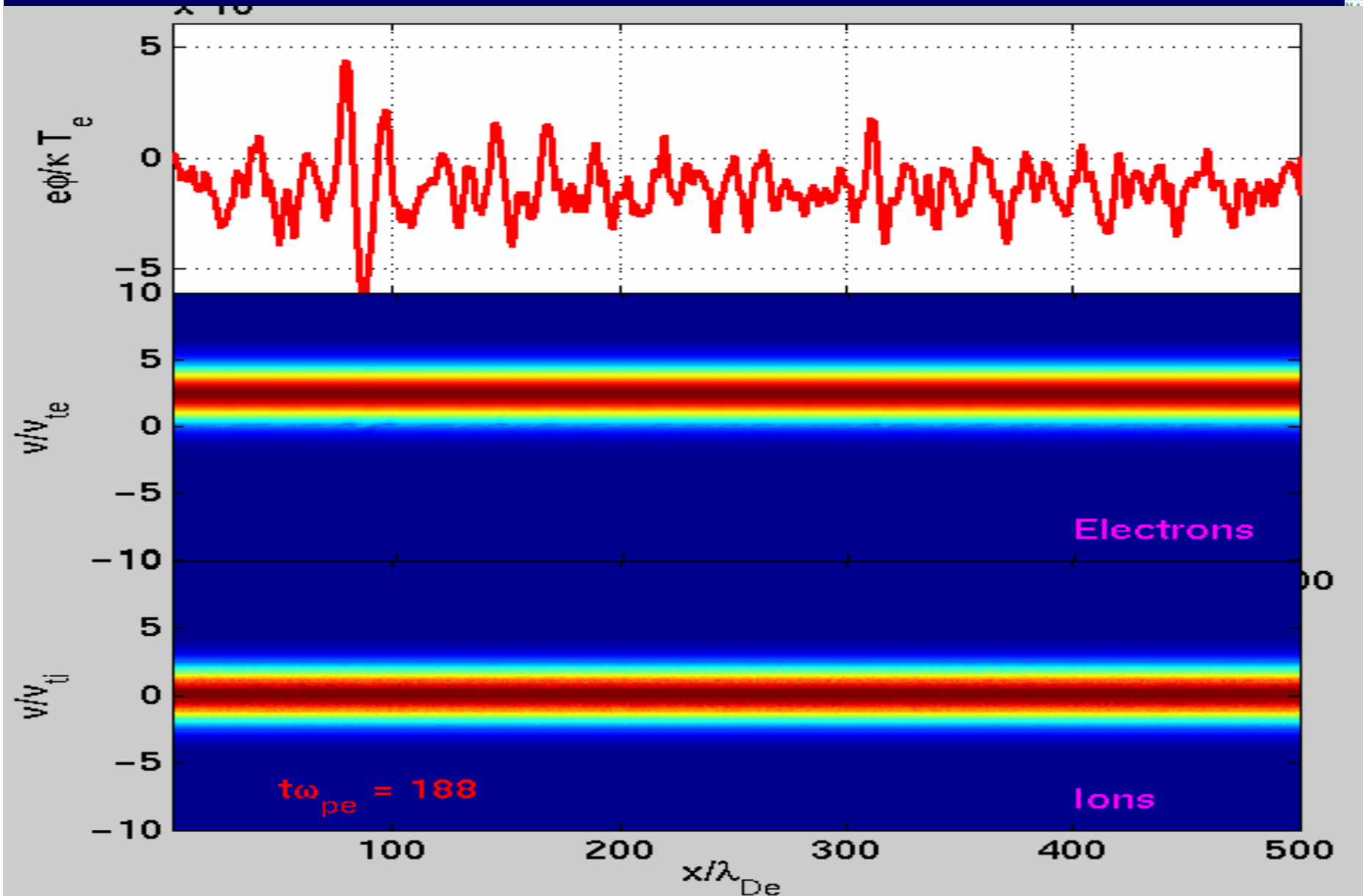
$$\frac{\partial f_{0e}}{\partial t} + \vec{v} \cdot \frac{\partial f_{0e}}{\partial \vec{r}} + \frac{e}{m_e} \vec{E} \cdot \frac{\partial f_{0e}}{\partial \vec{v}} = -\frac{e}{m_e} \left\langle \left(\delta \vec{E} + \vec{v} \times \delta \vec{B} \right) \cdot \frac{\partial \delta f_e}{\partial \vec{v}} \right\rangle$$

-> correlation of e/m fluctuations and plasma density /current fluctuations

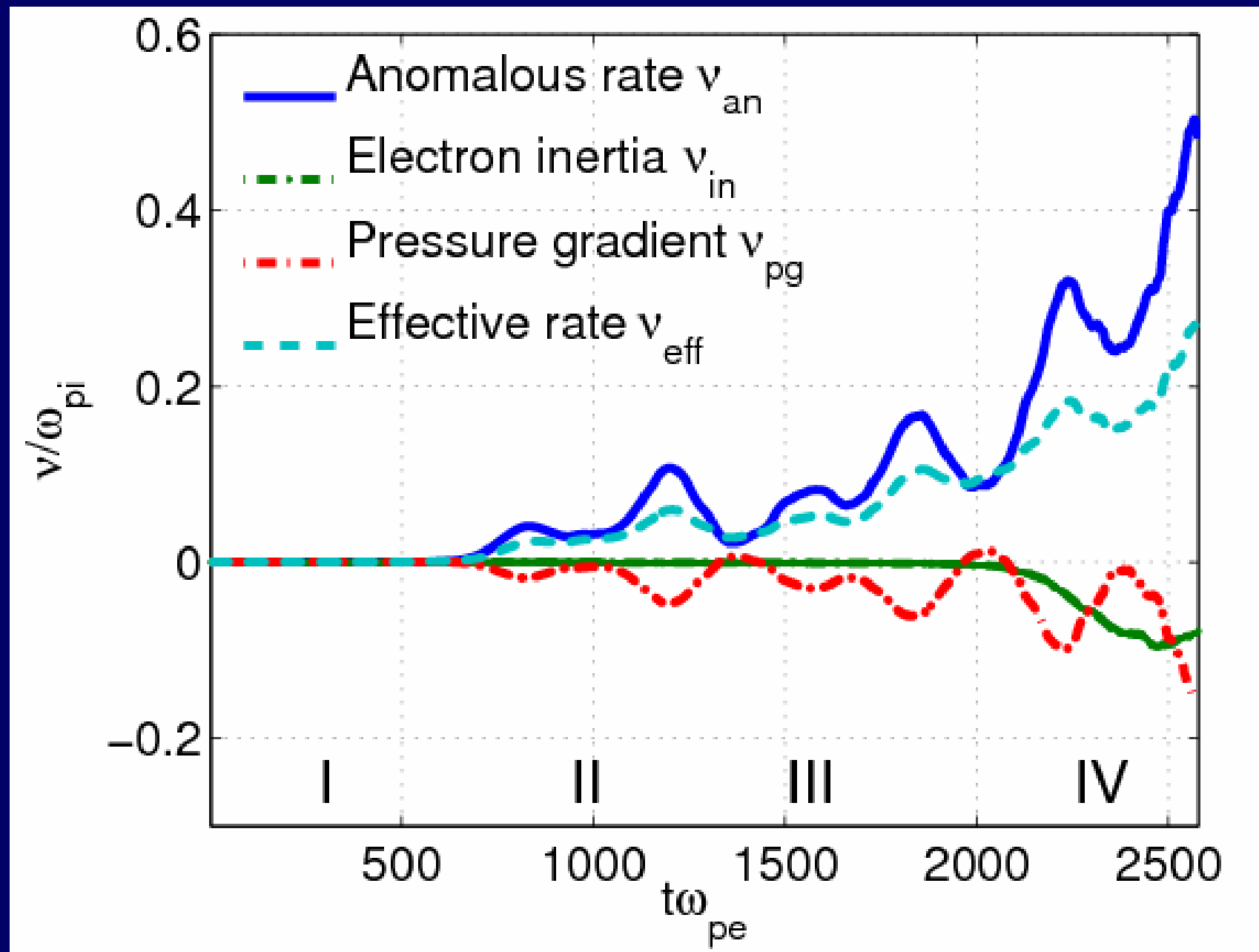
$$\left(\frac{d}{dt} n m_e v_{y,e} \right)_{eff} = \langle \delta E_y \delta \rho_e + \delta j_{z,e} \delta B_x - \delta j_{x,e} \delta B_z \rangle$$

-> These correlations can be taken either from theory, from observations or from numerical simulations !

Vlasov code kinetic simulation: evolution of potential, Fe and Fi



Effective „collision rates“



Dissipation in the solar atmosphere



Current dissipation, characterized as resistivity via an effective „collision frequency“ -> is dominated

$$\eta = \frac{\nu}{\epsilon_0 \omega_{pe}^2}$$

In the (lower) chromosphere:

by binary particle collision rate

[Spitzer, Härm & Braginski 1958-63]

$$\nu_{coll} \approx \frac{\omega_{pe}}{n \lambda_D^3}$$

In the corona:

by plasma turbulence as obtained by Vlasov code simulations for coronal conditions, Te~Ti etc:

[Büchner & Elkina 2006/2007]

$$\nu_c \approx \omega_{pi} / 2\pi$$

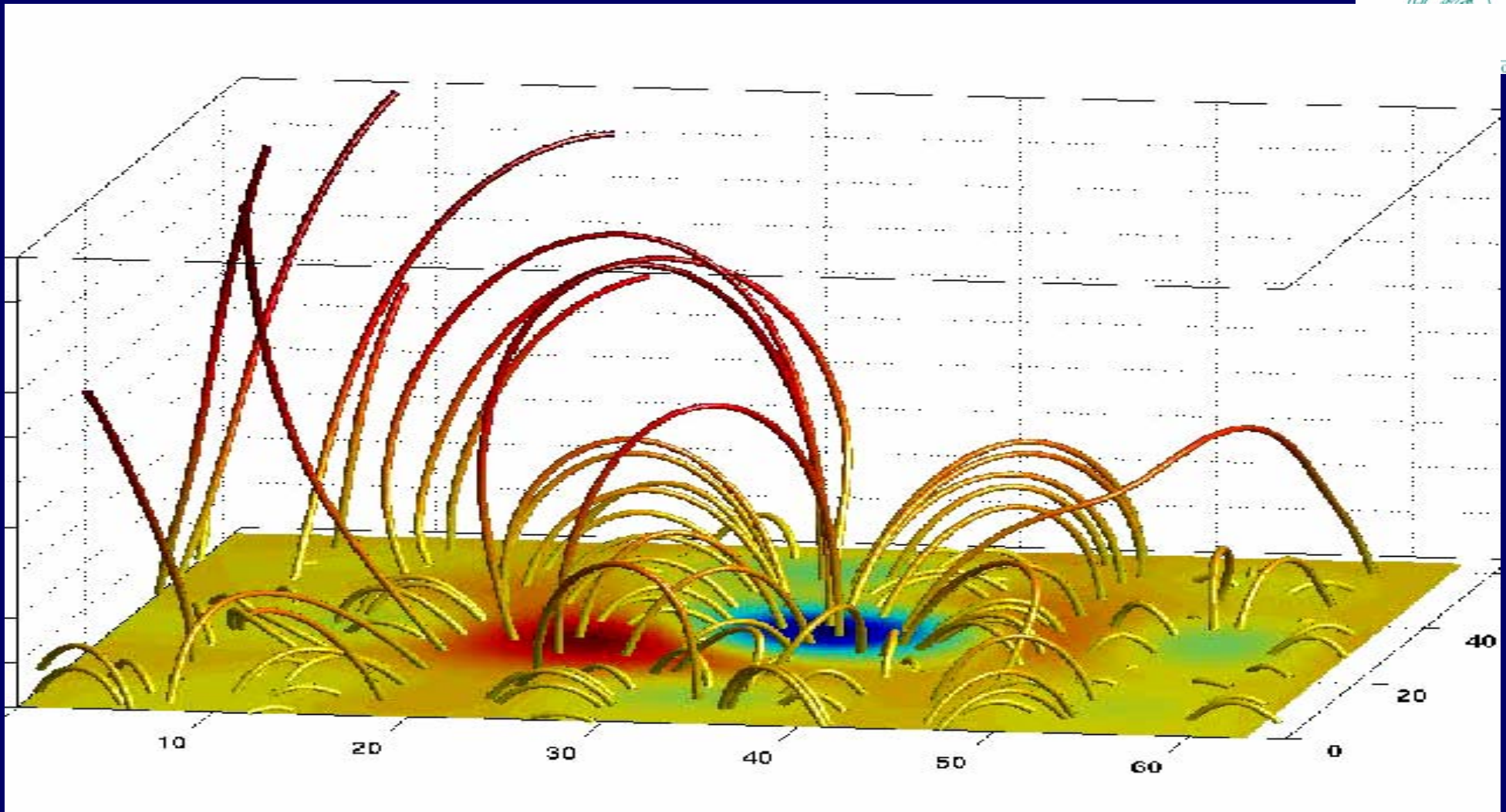
– for higher beta plasma -> 1D: IA double layers

– for lower beta plasma -> 2D: LH turbulence

$$\nu_c \approx \omega_{pi}$$

But: as threshold: a large current carrier drift velocity $j/ne > v_{te}$ (-> thin sheets!)

Resulting magnetic reconnection

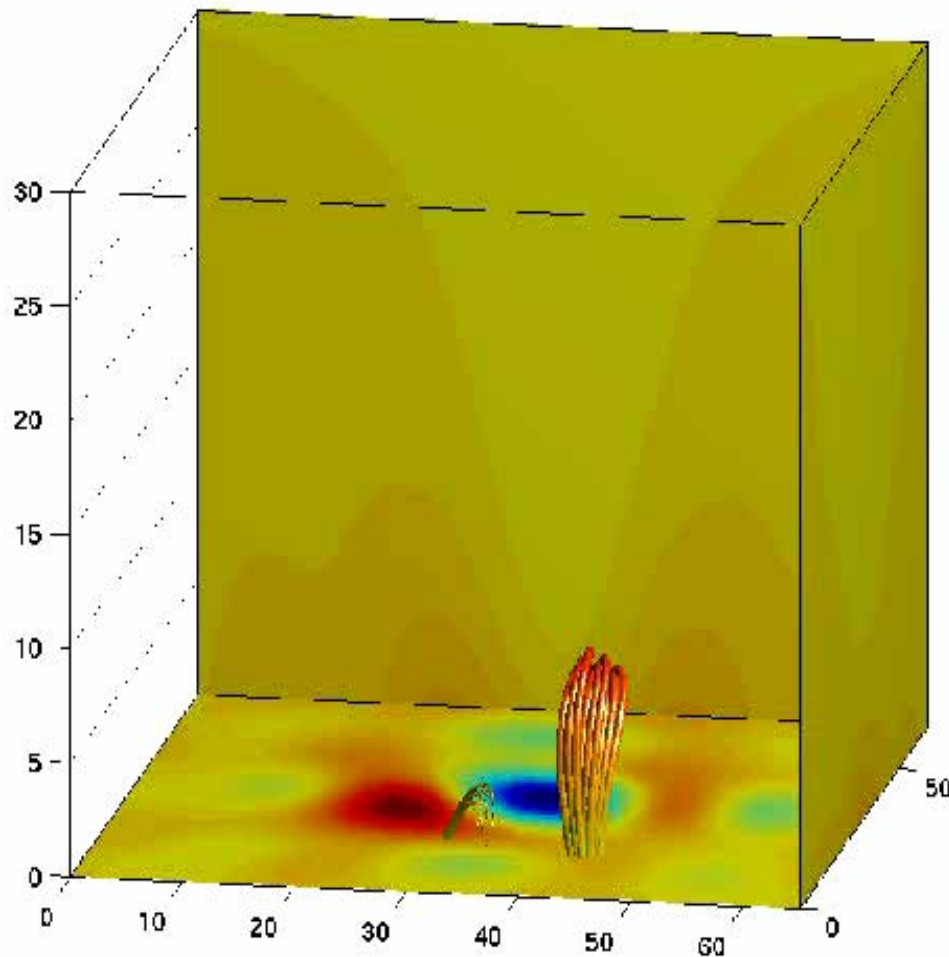


When currents caused by the photospheric plasma motion are dissipated -> reconnection

3D magnetic reconnection



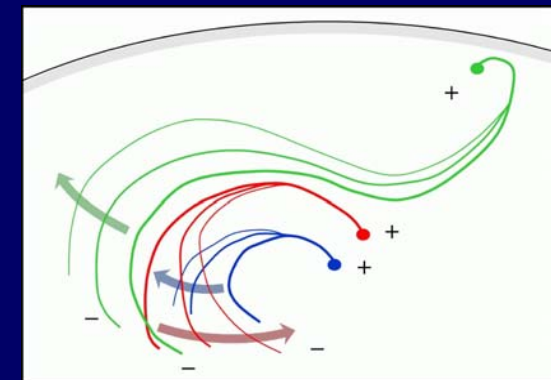
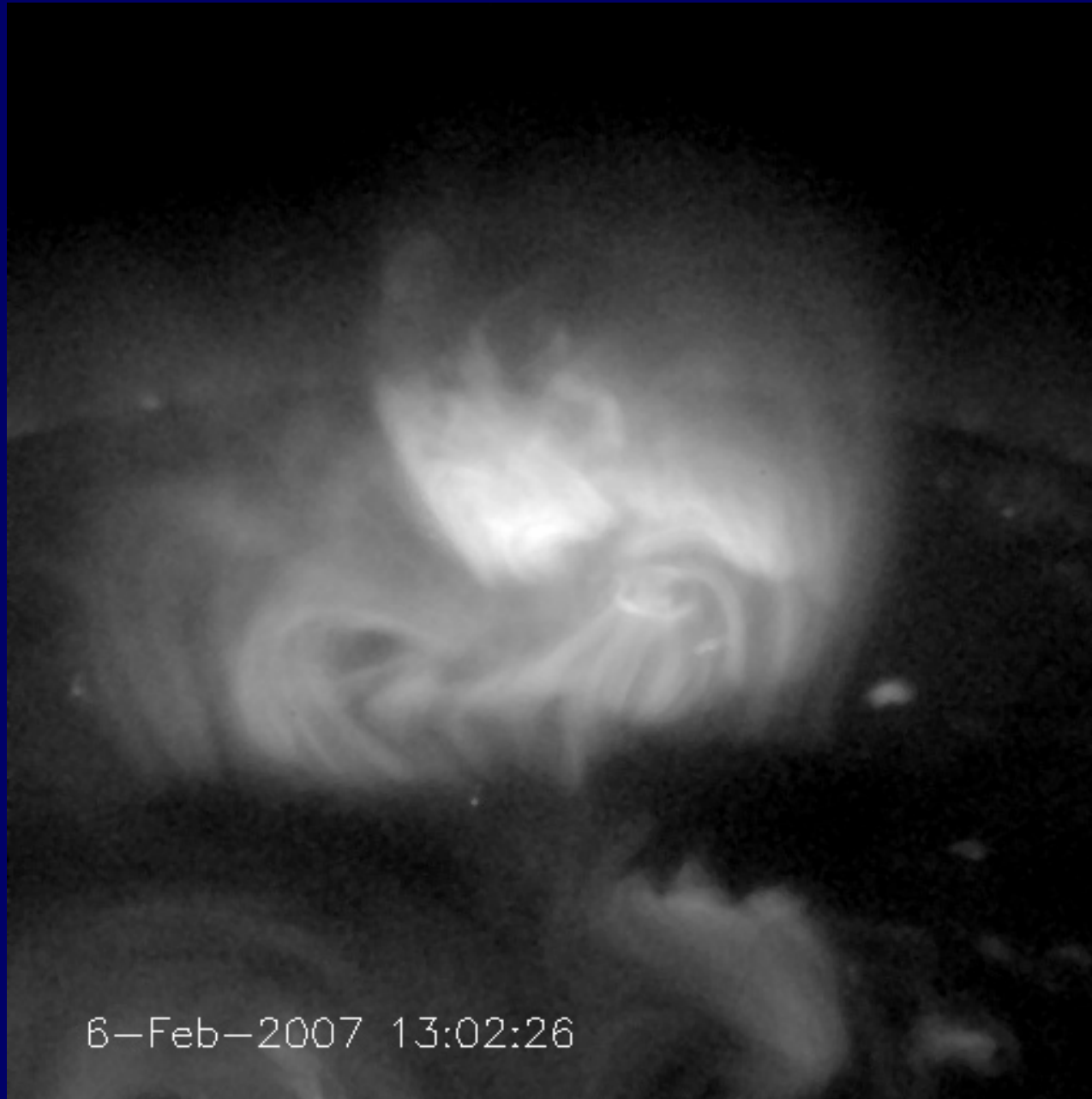
3D magnetic reconnection dynamics
[Büchner 2008]



Observation: Hinode 2007

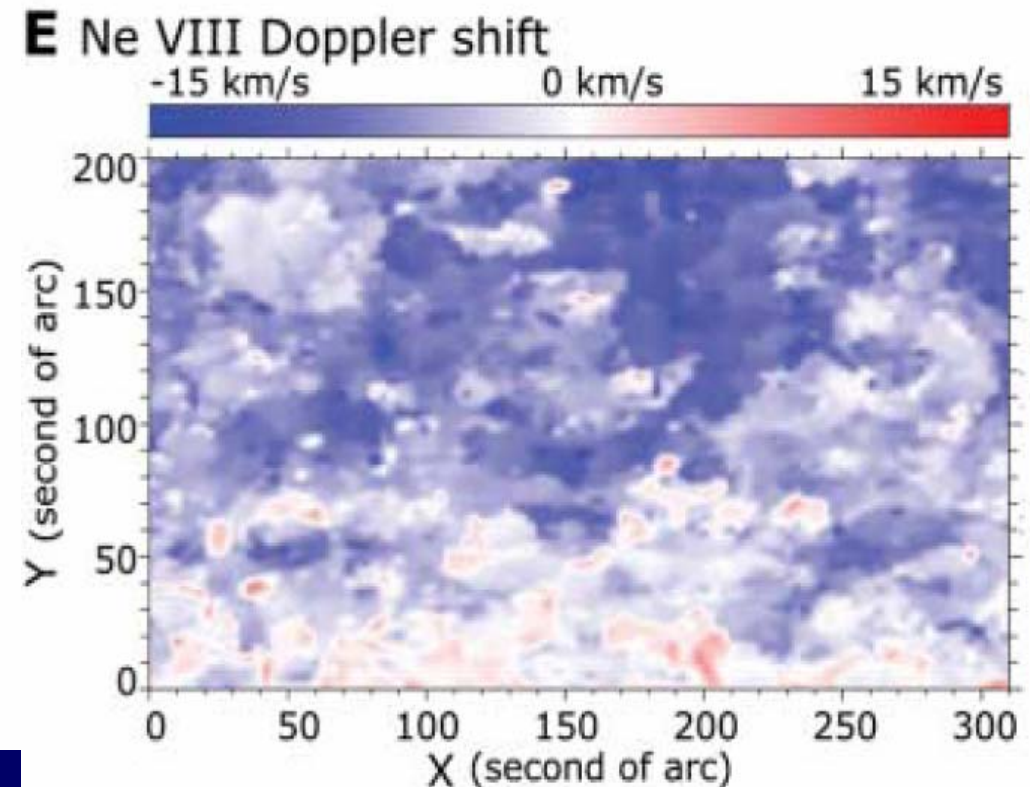
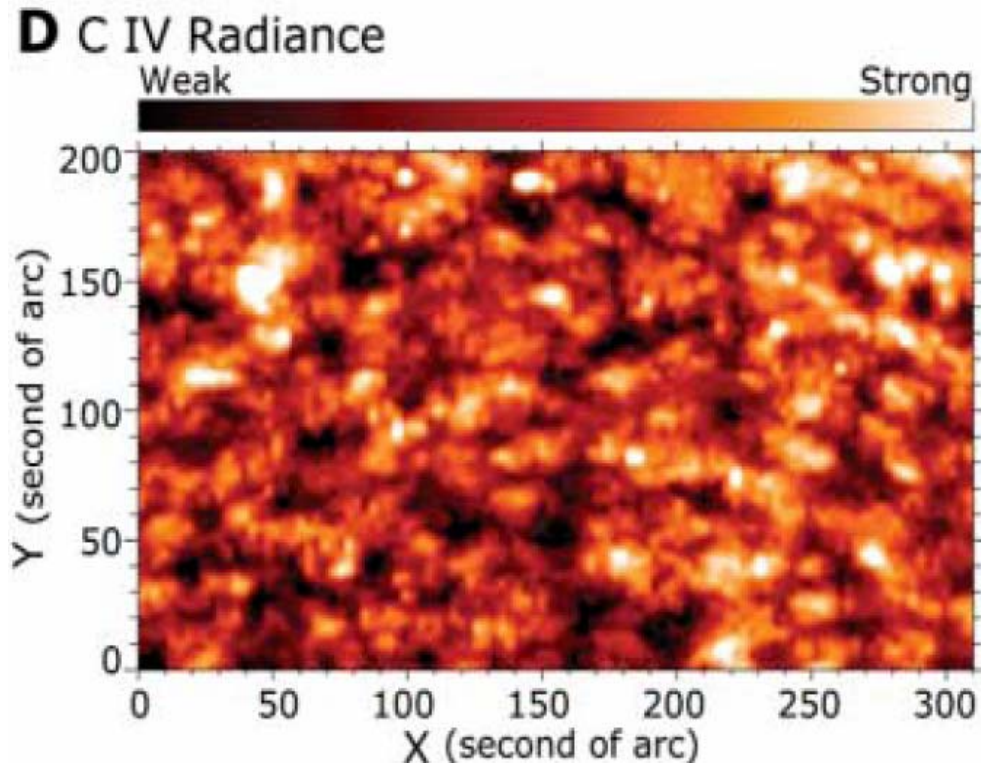


Japanese Space
Solar Telescope
mission (launch:
2006) ->
Observation of the
imaging X-ray
telescope
XRT



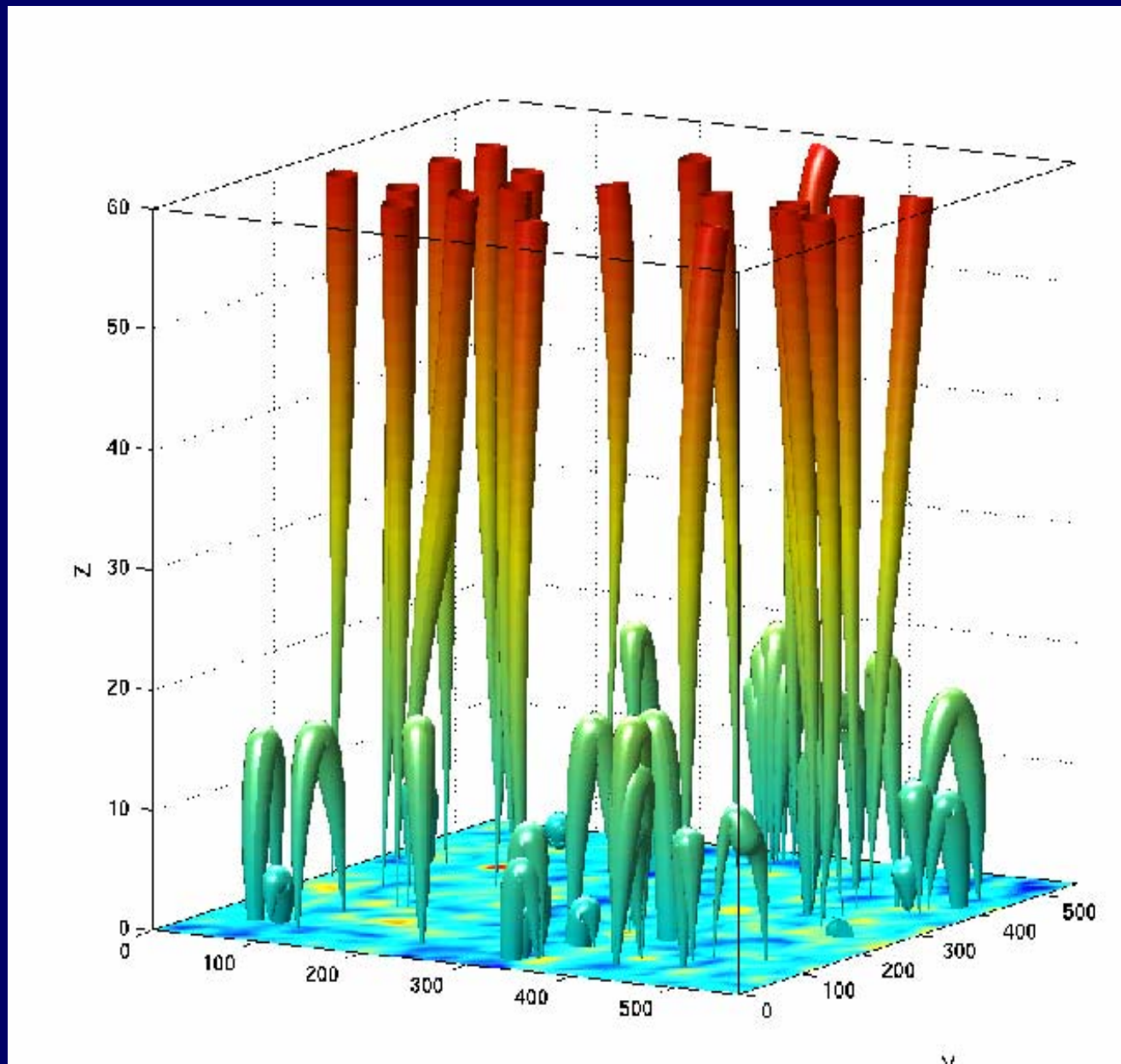
Sketch

1.) Solar wind from coronal holes



-> Observed radiance and velocities [Tu et al., 2005]

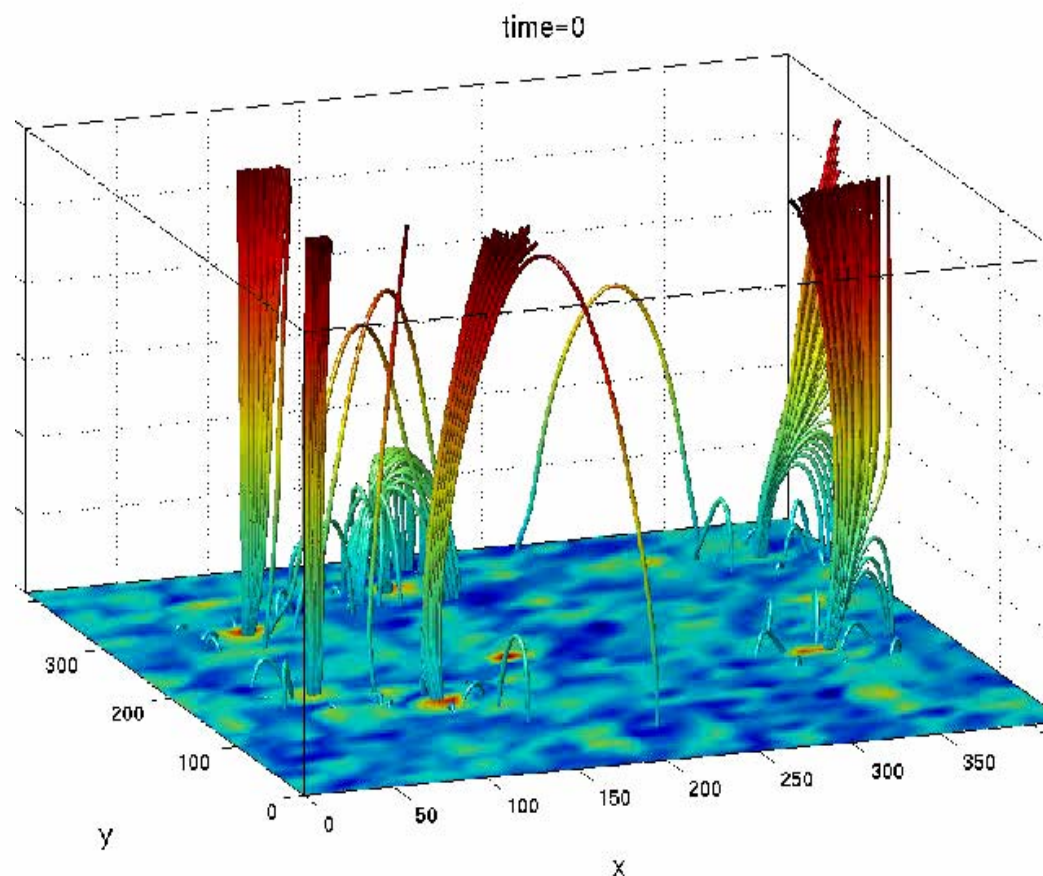
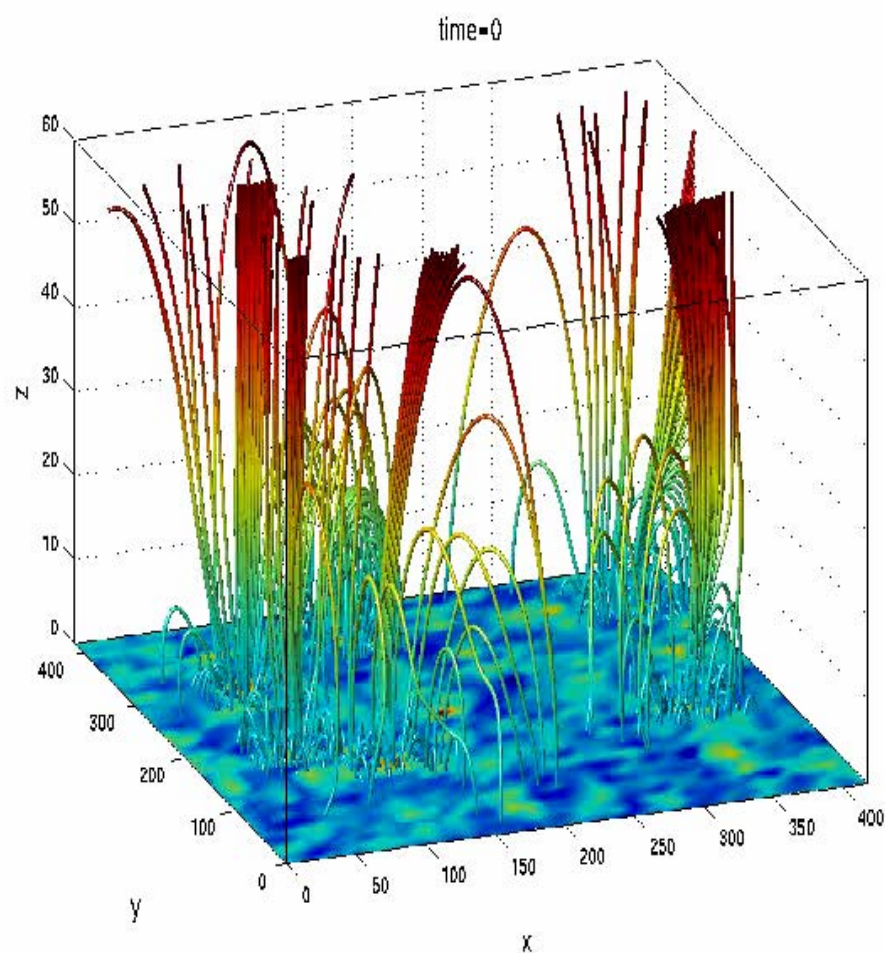
3D initial force free magnetic fields



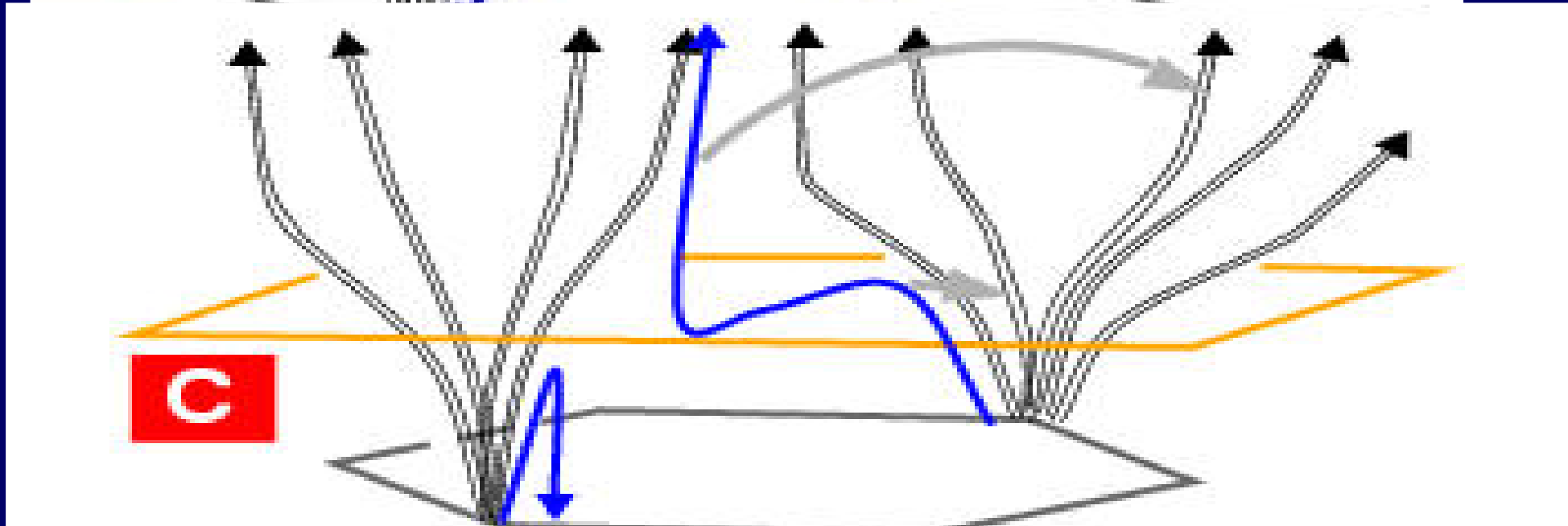
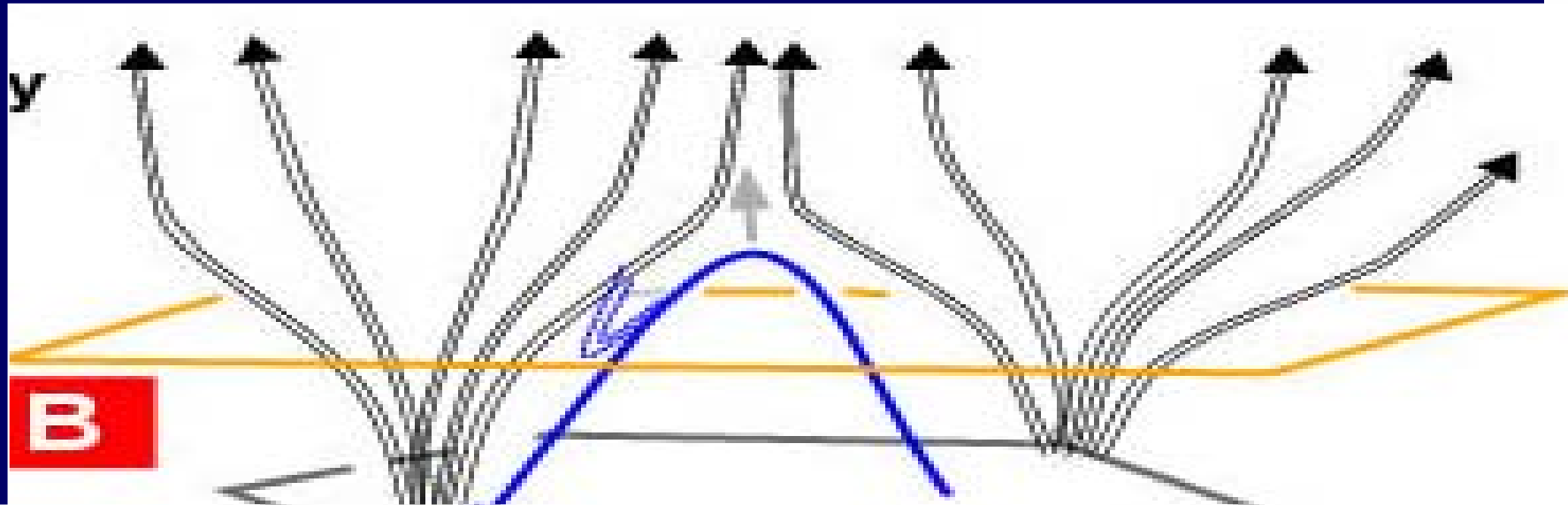
The magnetic field consists of funnels – open magnetic flux along the boundaries of the coronal hole and low rising loops inside and outside the hole

[Büchner, 2010]

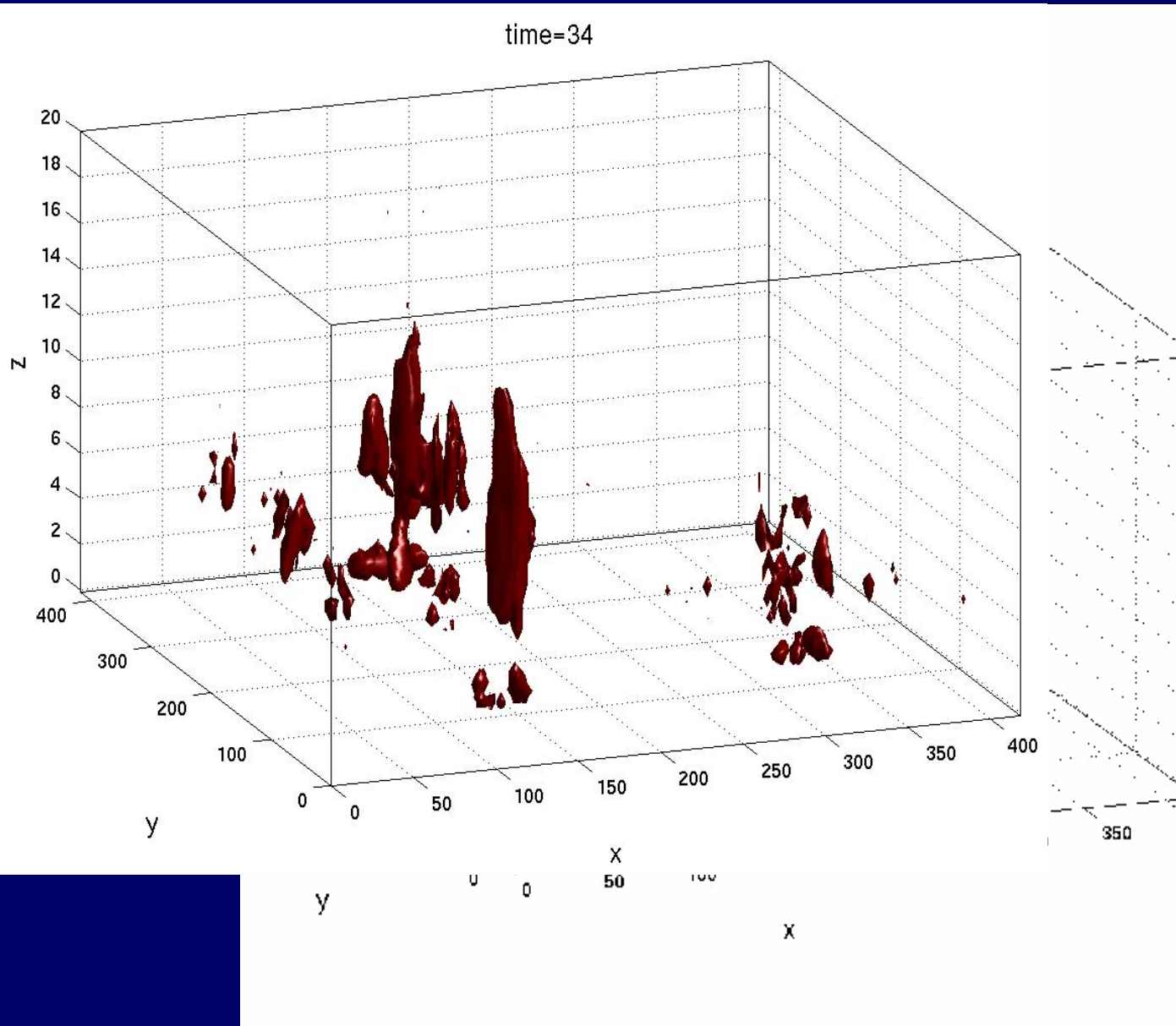
Resulting B field evolution



Reconnection - cartoon



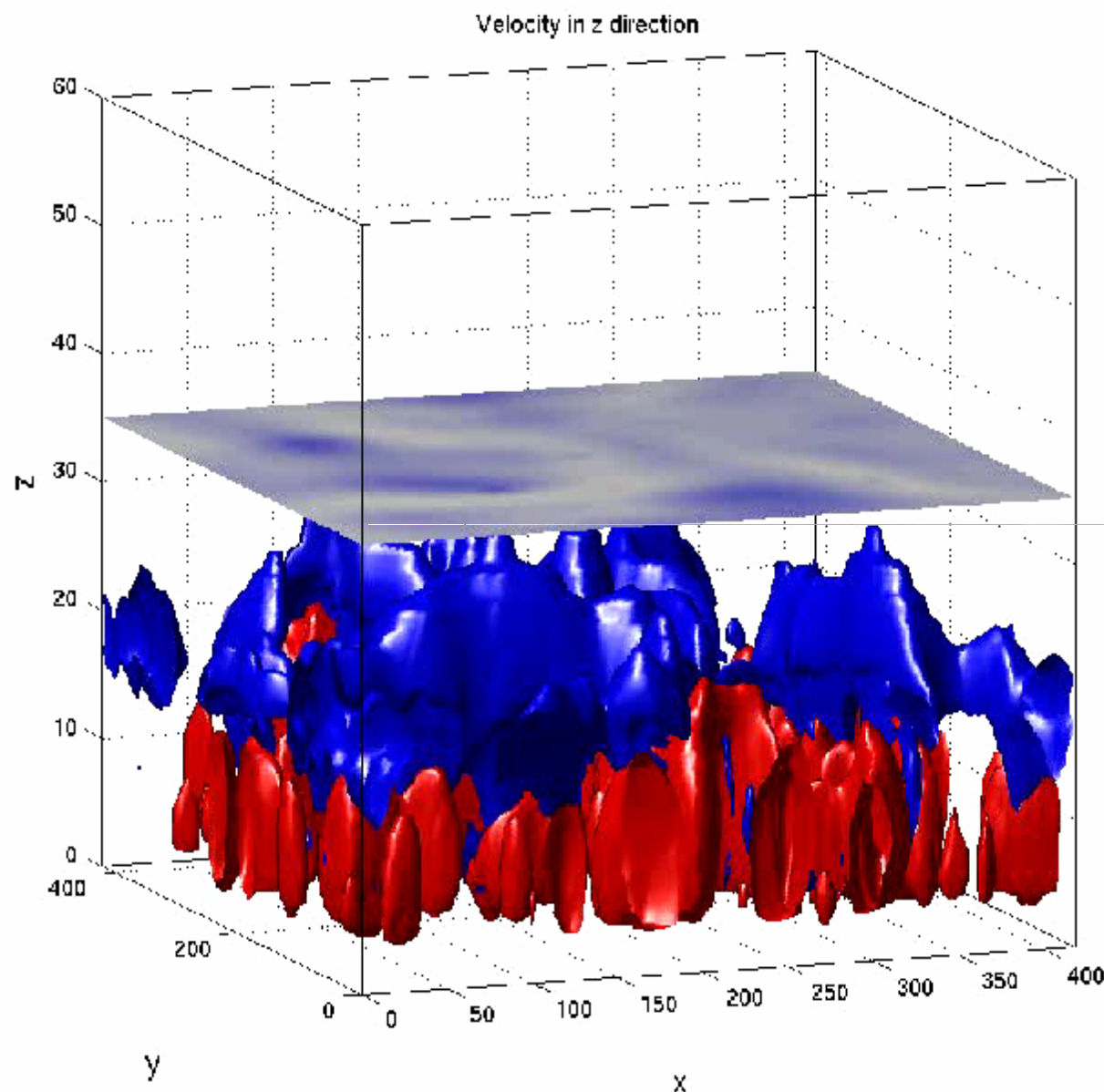
Reconnection: Epar field



3D
reconnection is
characterized
best by the
parallel electric
field, shown to
the left

-> Epar and,
therefore, the
3D
reconnection
rate is highest
between 5 and
10 Mm above
the
photosphere

Epar -> Plasma acceleration



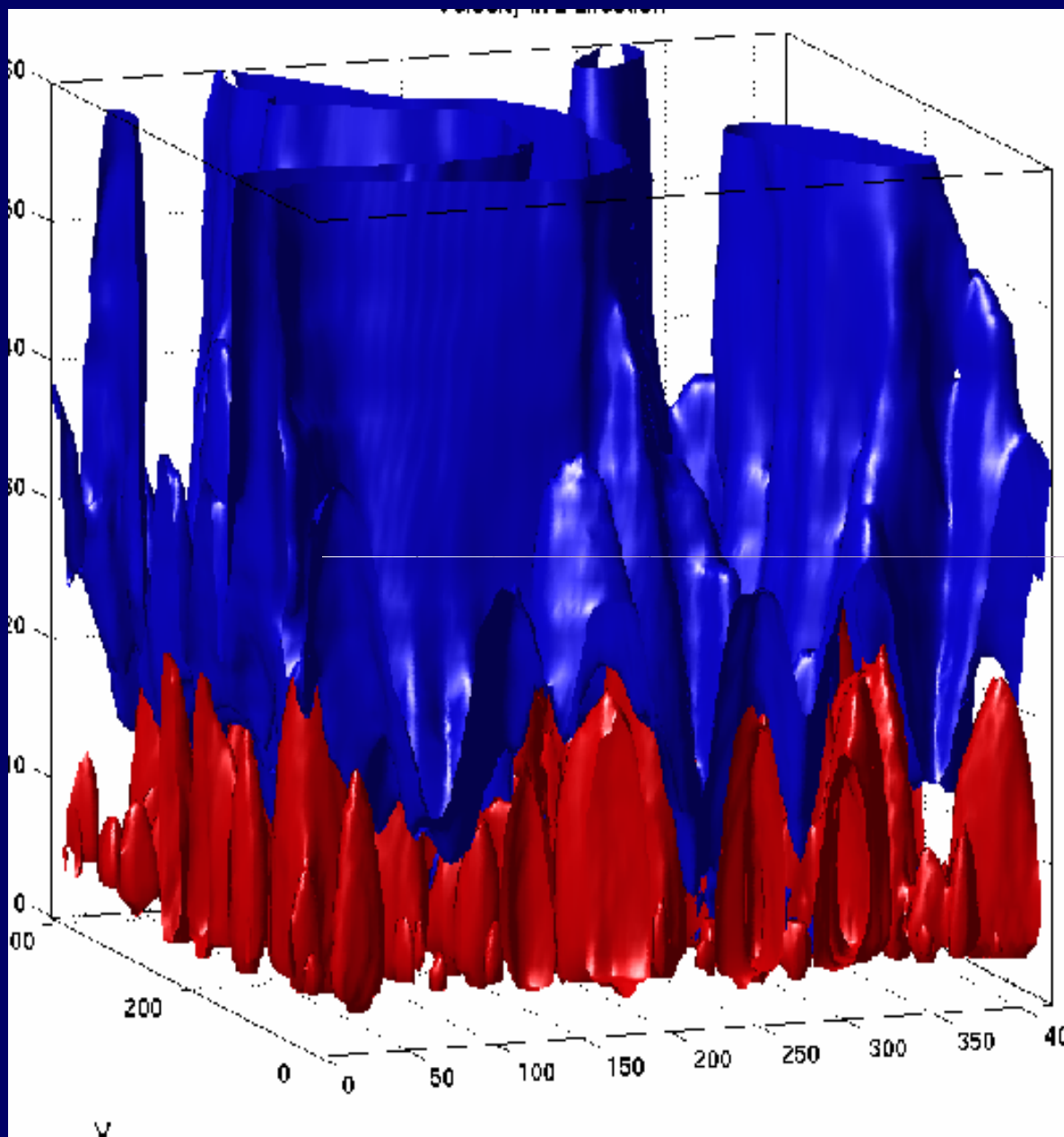
The isosurfaces
 $V_z=10$ km/s

- blue: upward
- red: downward

demonstrate:

3D reconnection electric
fields accelerate plasma
both upward and
downward

Accelerated solar wind plasma



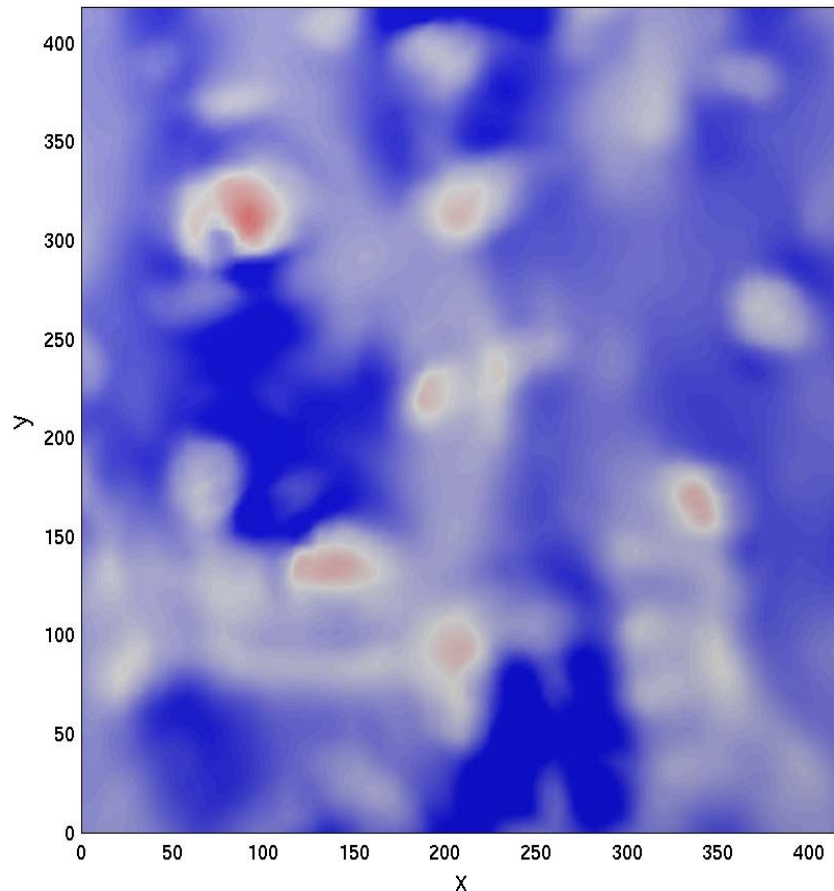
These stationary upward (blue) and downward (red) accelerated plasma flows – velocity isosurfaces 10 km/s show:

- the upward directed acceleration takes place mainly above 5 Mm and
- higher up the accelerated flows fill almost the whole coronal hole area

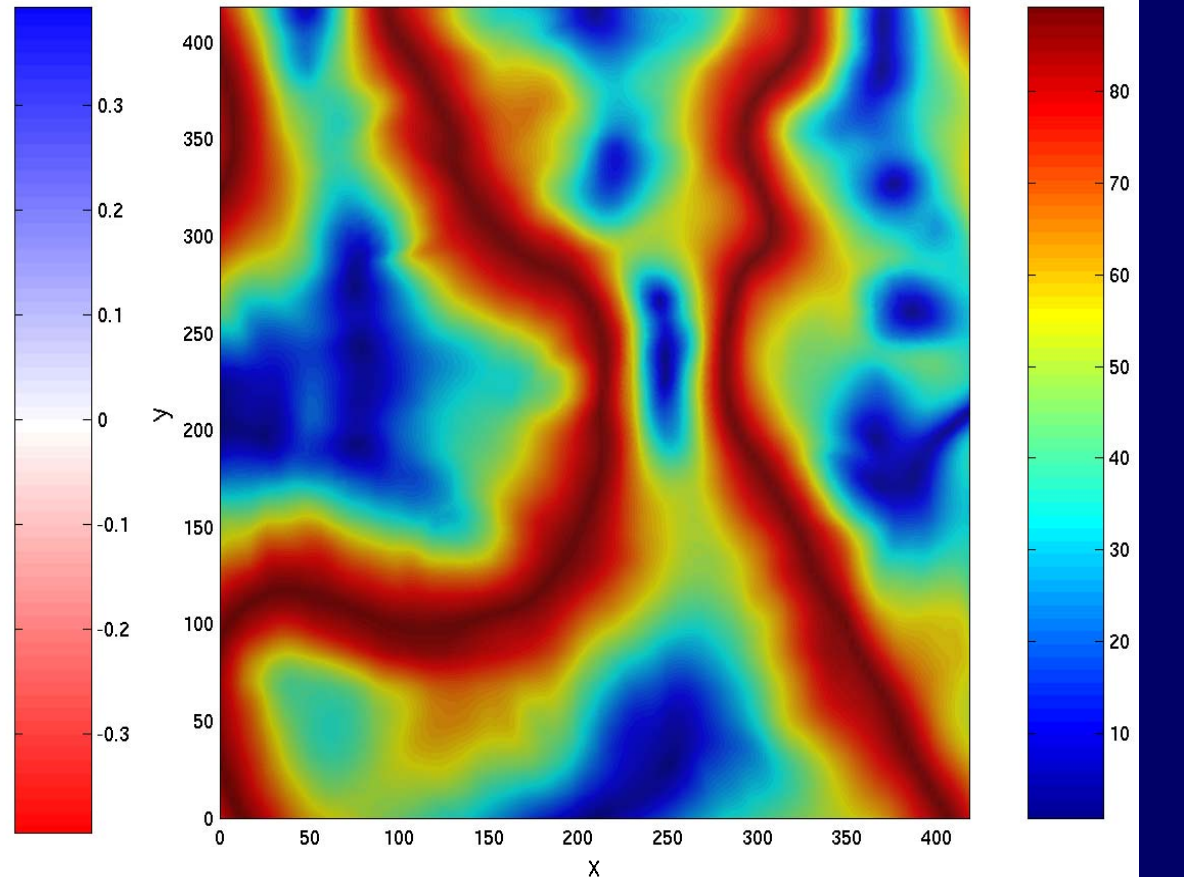
Simulated V_z and B -inclination



Velocity in z direction



Inclination at $h_z = 35$

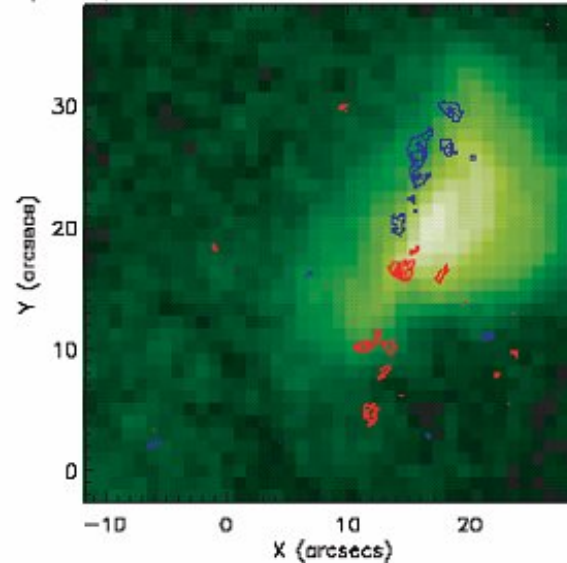
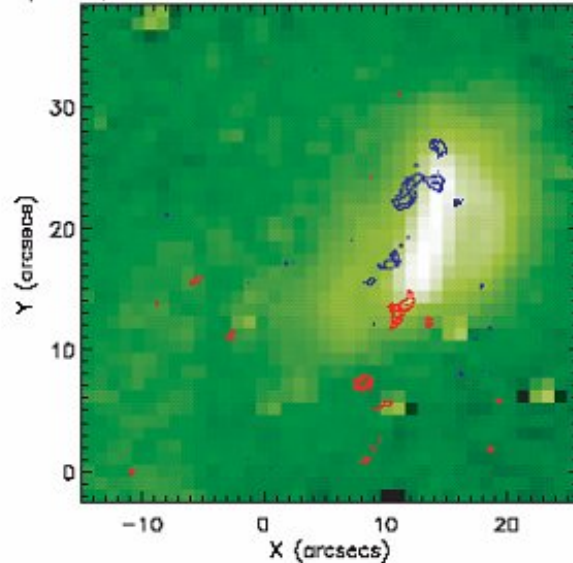


Comparison V_z (left) and B -inclination (right) at $H=20$ Mm $\rightarrow$ clear correlation: downward and stagnant flows on closed field lines (small inclination) and upward flows on open field lines (large inclination)

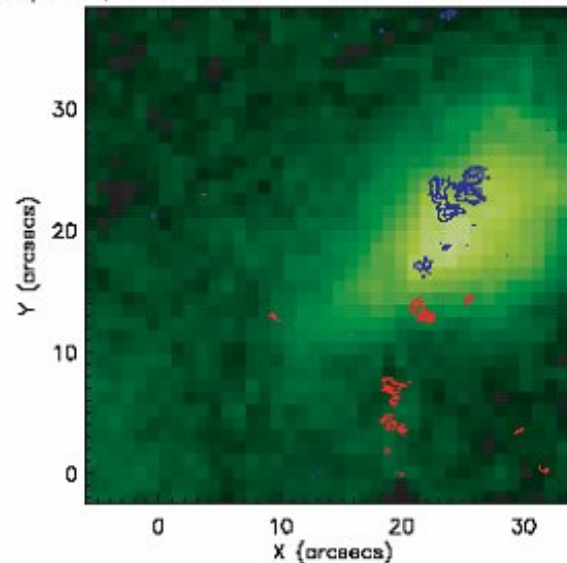
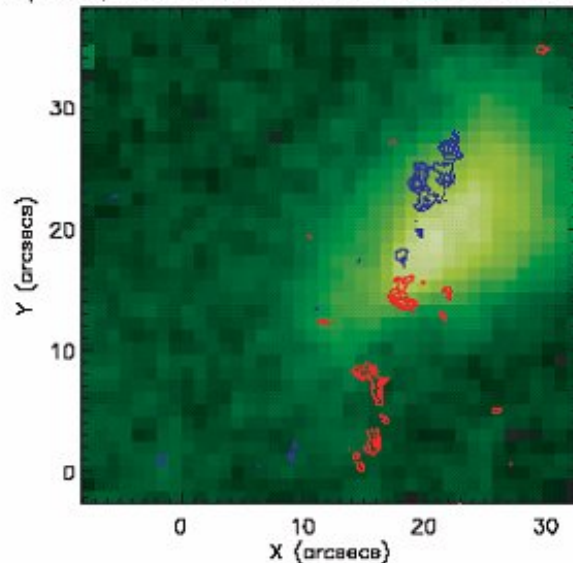
2.) X-ray Bright Point (Hinode)



JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:00:44.397 UT JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:21:42.816 UT



JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:43:03.241 UT JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:56:50.014 UT

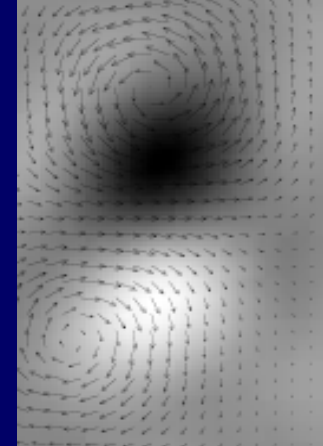
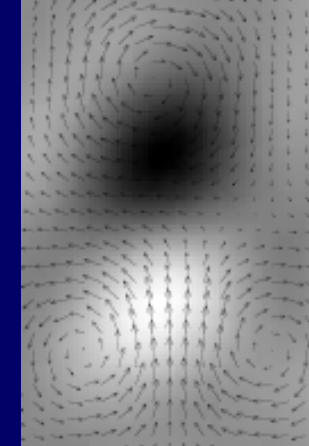
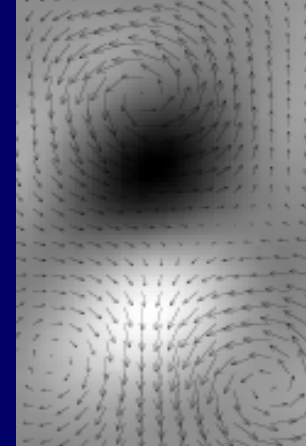
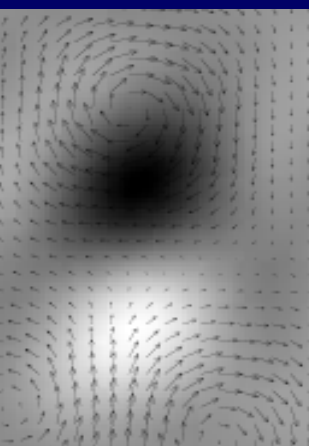
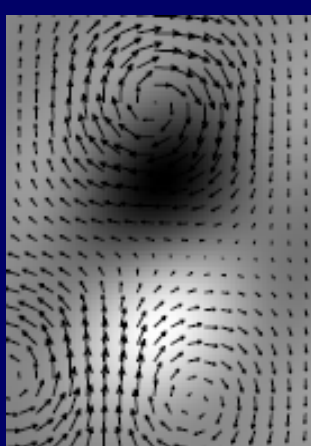
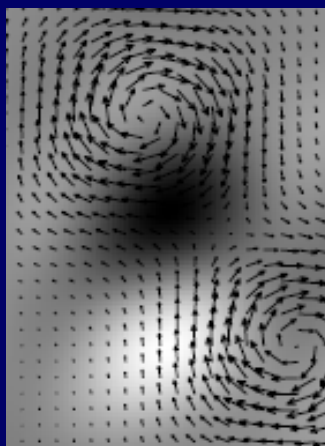
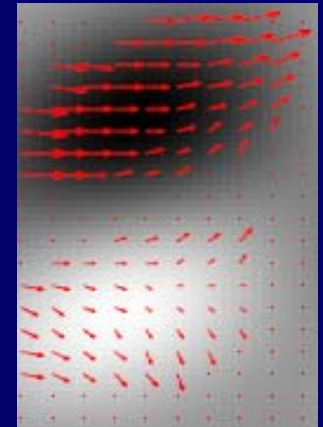
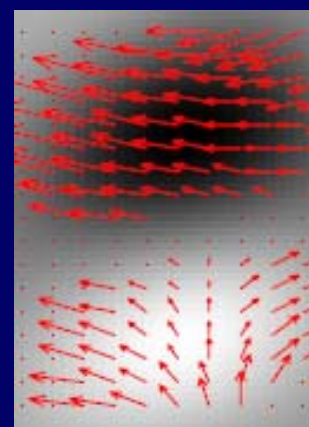
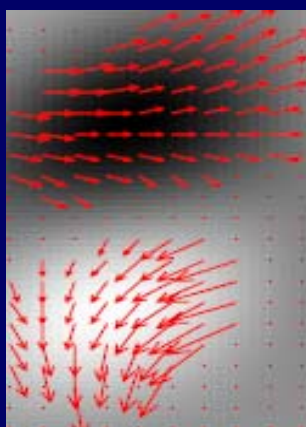
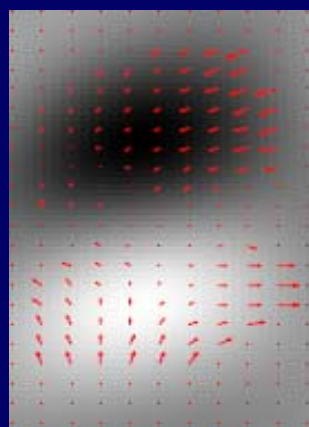
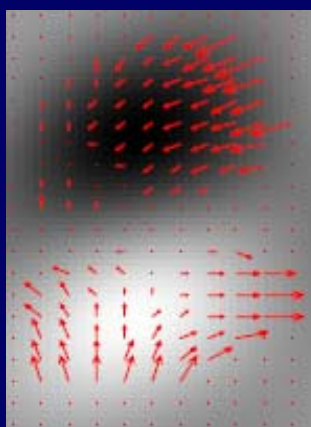
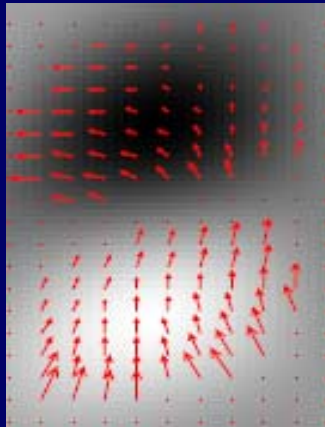


These four
X-ray images
were observed
by Hinode XRT
between
23:00 UT
and
24:00 UT
on December
12, 2006

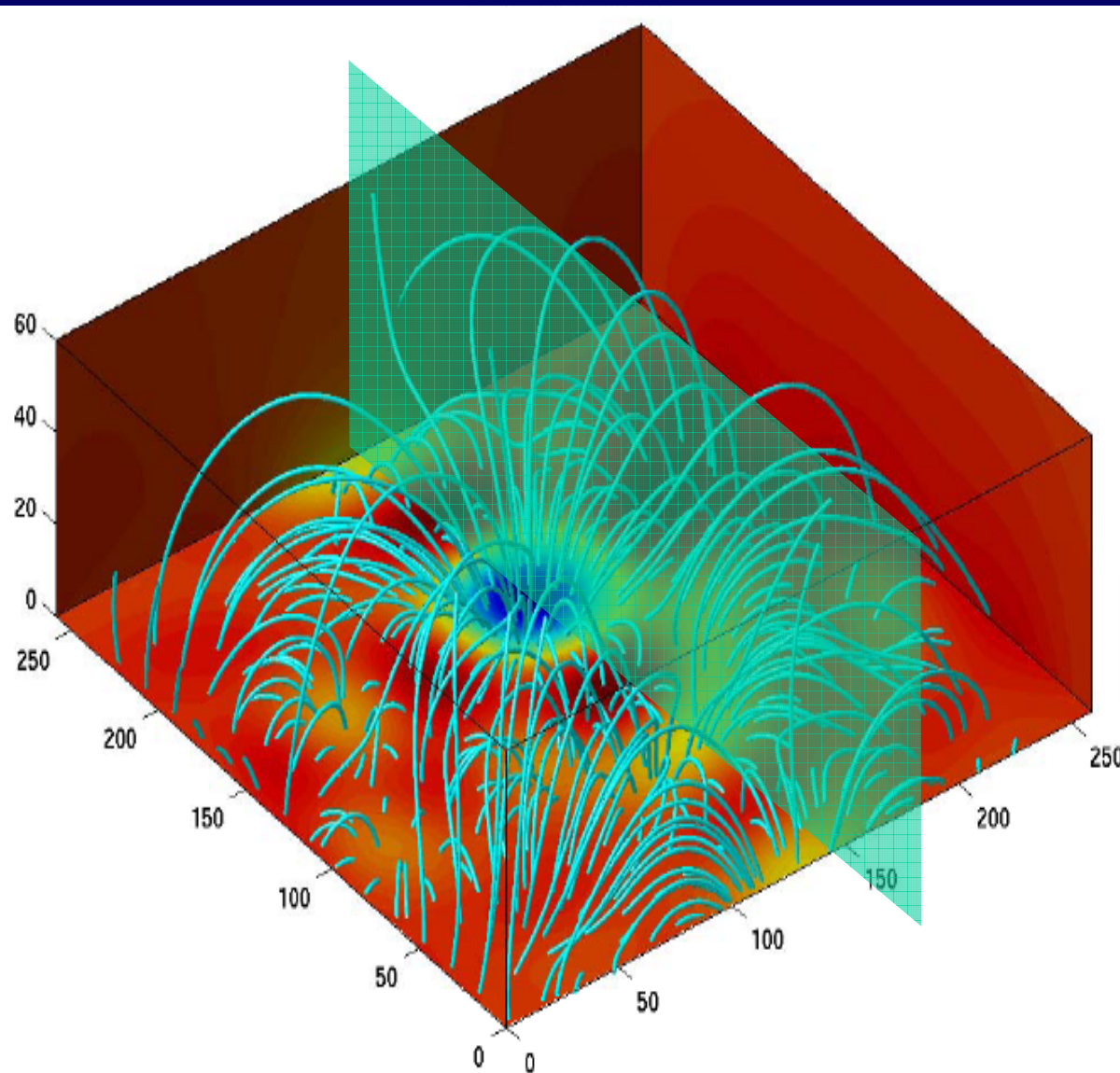
Energy input by various types of photospheric plasma motion



23:02 – 23:07UT	23:07 – 23:12UT	23:12 – 23:17UT	23:17 – 23:22UT	23:22 – 23:27UT	23:27 – 23:32UT
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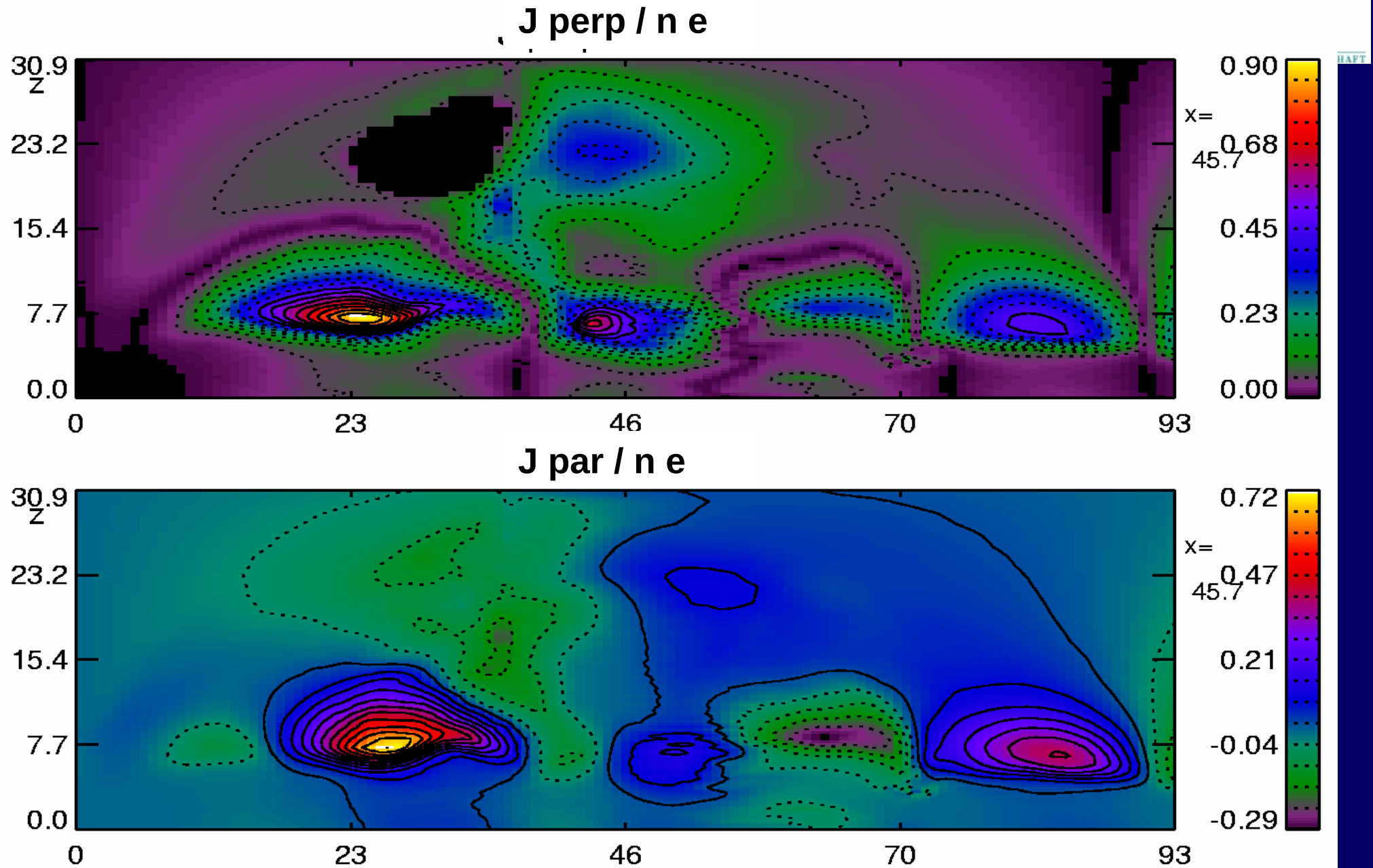


Simulation results

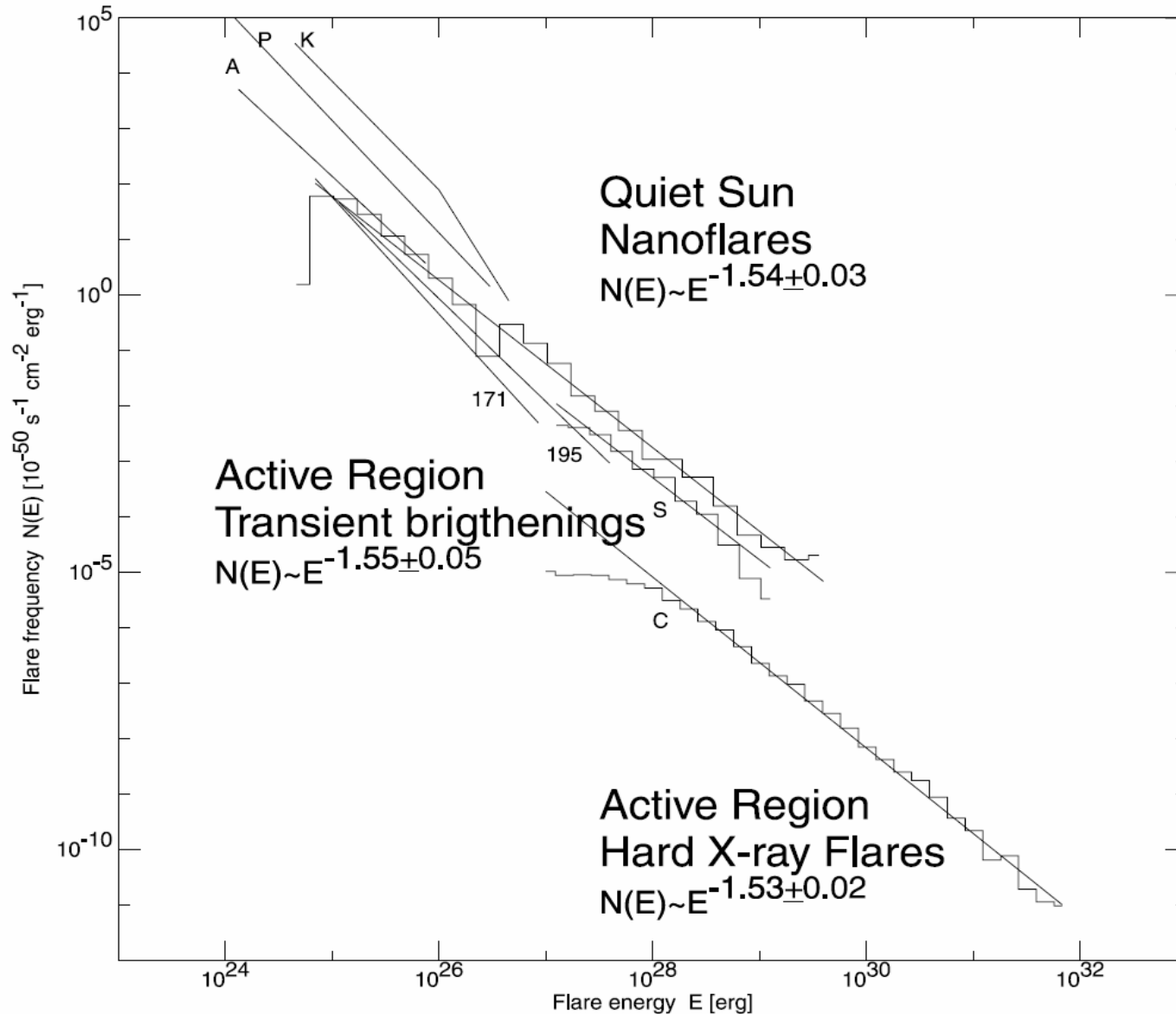


-> not any more
force-free, i.e. J_{perp}
-> The basic features
are robust, do not
depend on the
details of the
photospheric
plasma motion.
Vertical cut:
plane of diagnostics
(through the
main polarities)

Generation of J_{par} and J_{perp}



Energy release at many scales

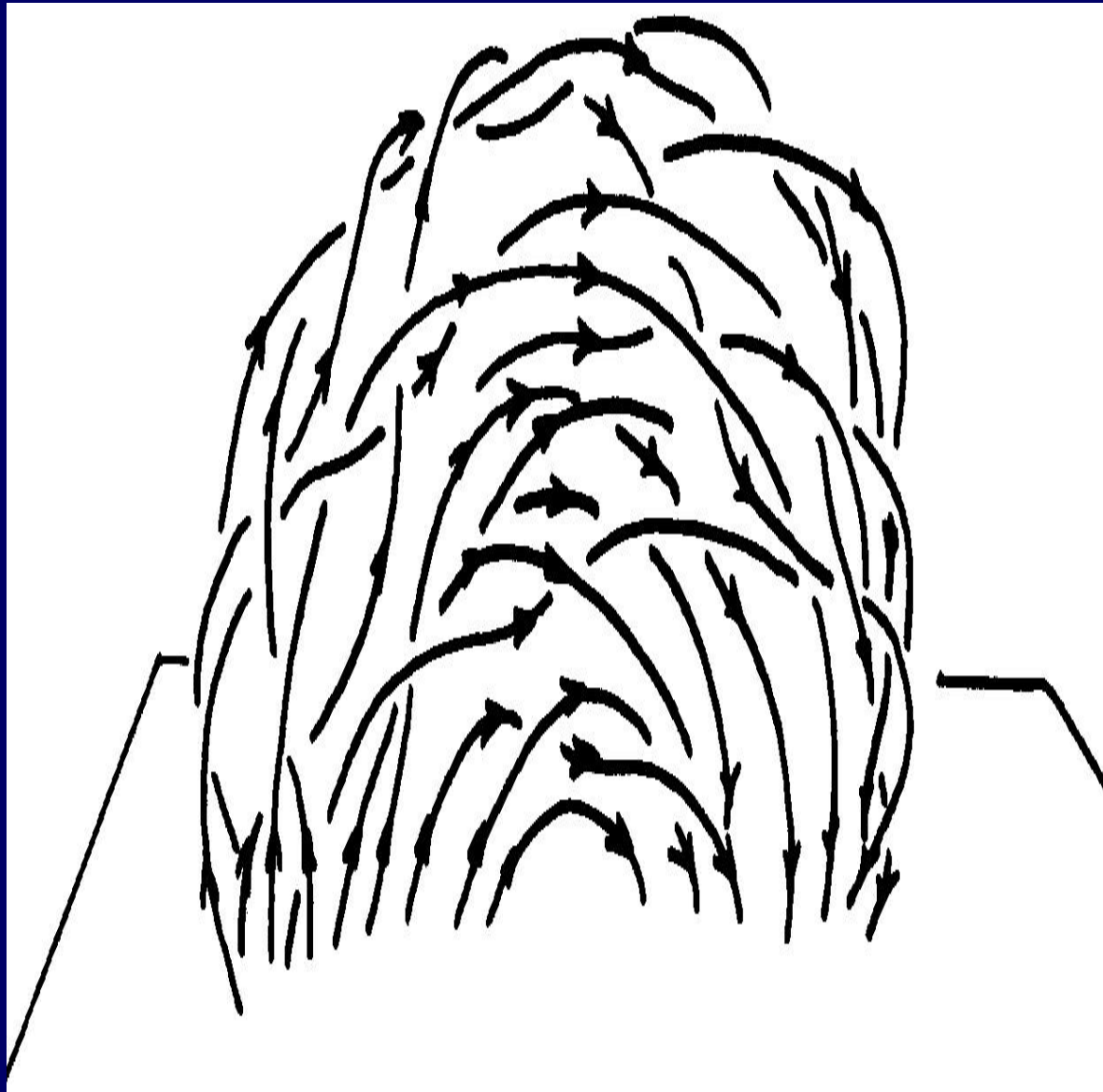


$N(E)$ is the probability of an event in dependence on its energy $dE \, dA \, dt$

The total number of events is
The integral over $N(E) \, dE \, dA \, dt$.

[Aschwanden and Parnell 2008]

Heating by nanoflares



Footpoint motion in
the photosphere

-> tangential
discontinuities

-> small scale
reconnection

-> '*Nano-Flares*'

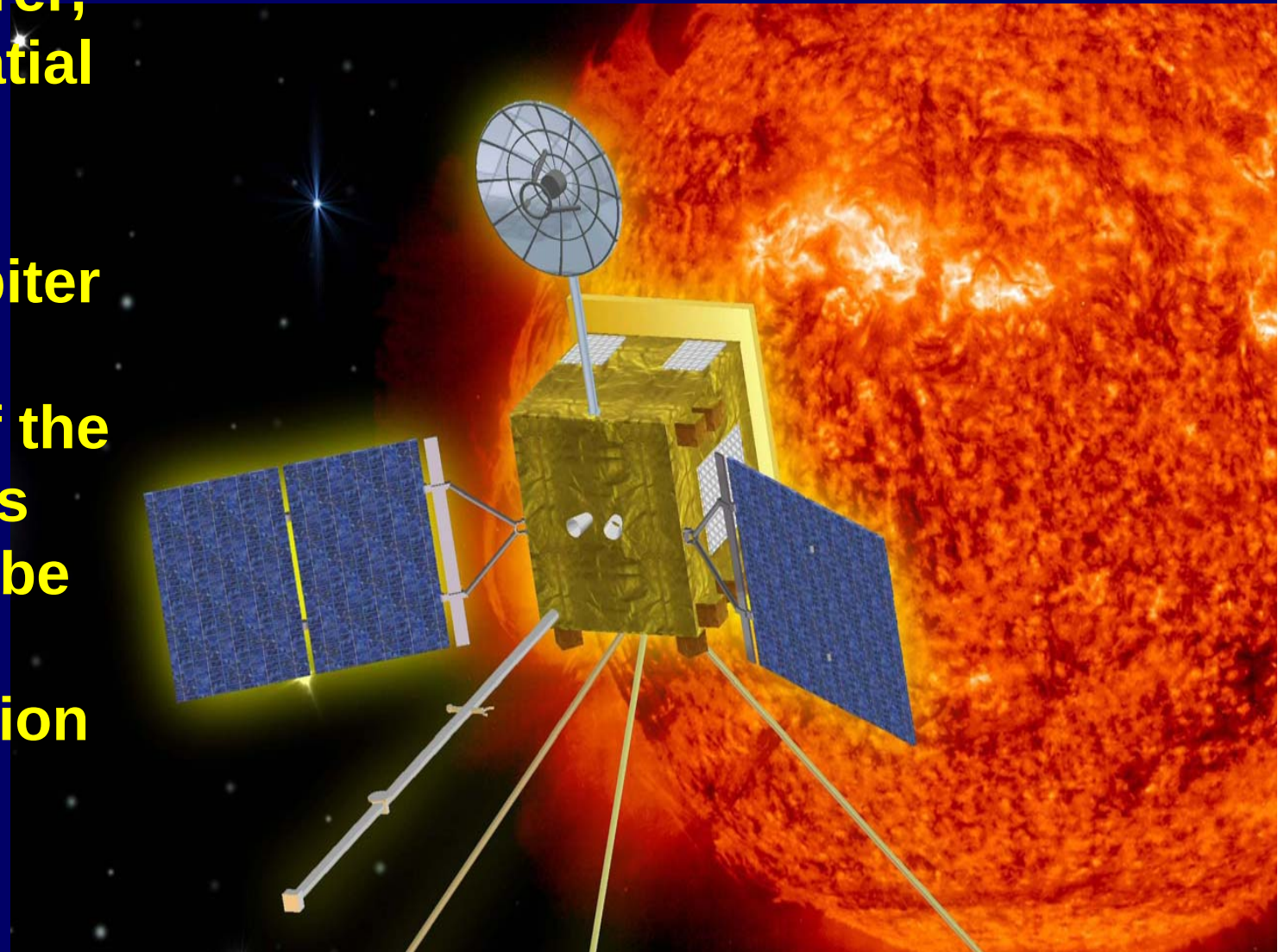
$E \sim 10^{24}$ ergs, $t \sim 1$ s

[Parker, 1988]

**But: Not observable
at Sun in principle !!**

Next steps:

- 2010 „SDO“ (Solar Dynamics Explorer, NASA): High spatial and temporal resolution
- 2017: Solar Orbiter (ESA):
Investigation of the Solar poles
- 2018: Solar Probe Plus:
In situ investigation of the solar atmosphere



Summary



- The heating and dynamics of the solar atmosphere is mainly due to its magnetic coupling to the photosphere
- The energy is initially transferred by the interacting plasma and neutral gas in the chromosphere
- Complex solar B field + photo-/chromospheric plasma motion -> Large scale 3D current systems
- Currents reach small (plasma) scales -> their dissipation is microscopic, to be described kinetically
- Combining RMHD and kinetic physics one can model
 - solar wind acceleration
 - coronal heating (bright points!)
 - heating by nanoflares